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THz studies with SR and FEL radiation

- The THz spectral range
- THz spectroscopy with Synchrotron Radiation
- THz non-linear studies



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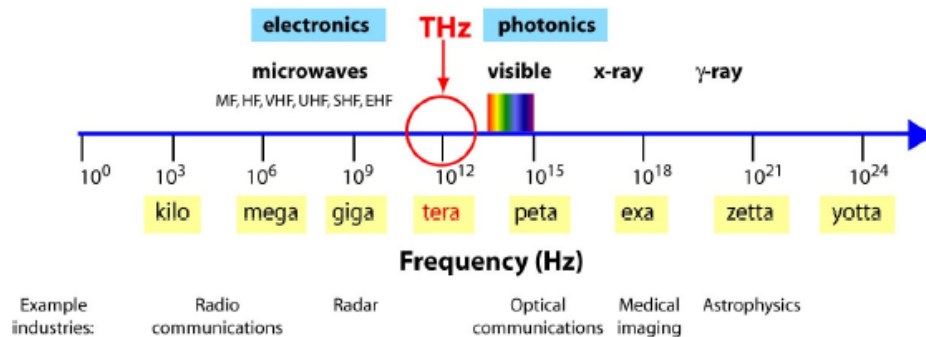
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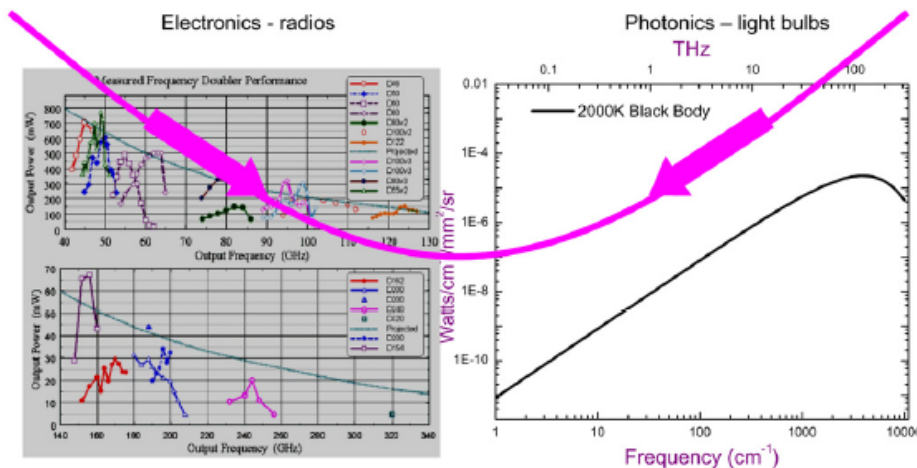
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The THz gap



$$1 \text{ THz} \sim 1 \text{ ps} \sim 300 \mu\text{m} \sim 33 \text{ cm}^{-1} \sim 4.1 \text{ meV} \sim 47.6^\circ\text{K}$$

Figure 1. Schematic of the electromagnetic spectrum showing that THz light lies between electronics and photonics.



G. Williams, Rep. Prog. Phys. 2006

Free electrons (Drude)

Lattice vibrations (phonons)

Rotational modes

Librational/Torsional modes

Intermolecular modes

Hydrogen bonds

Van der Waals interactions

Electro-magnons



Properties of THz light

- **Non-ionizing**

safe use on living people/animals, non-destructive for biological samples

- **Highly penetrating**

sees through many materials, as packaging, clothing, walls

- **Chemical specificity**

distinguishes between different plastics, drugs, explosives

- **High contrast**

between strong (metals, water) and weak (plastics, tissues) absorbers

- **High-speed communications**

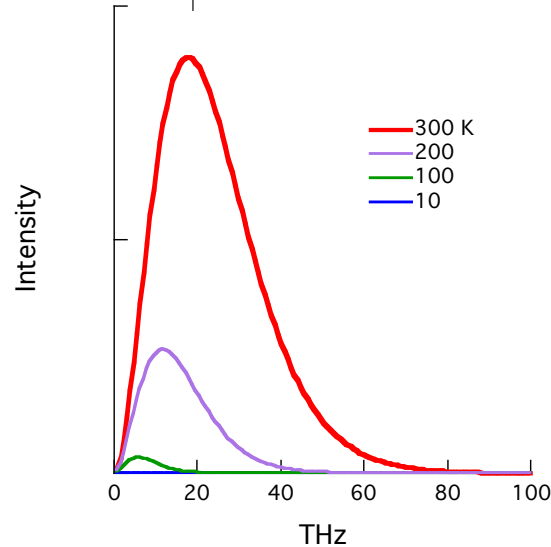
1000 times faster than GHz

...the main drawback is spatial resolution ~ mm



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THz light and blackbody radiation



Planck's law

$$I(\nu, T) = \frac{2h\nu^3}{c^2} \frac{1}{e^{\frac{h\nu}{kT}} - 1}$$

Wien's law

$$\lambda_{\max} = b/T \quad b = 2.897 \cdot 10^{-3} \text{ K m}$$
$$\nu_{\max} = 0.1035 T \quad \text{THz K}^{-1}$$



All objects at room temperature are THz radiation emitters



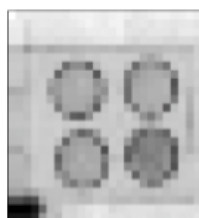
THz modes are normally populated at room T

Technological applications

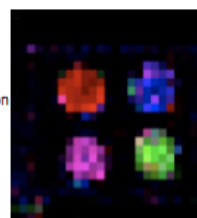
Pharmaceutical



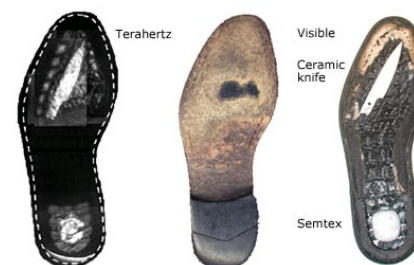
Fig. 8. Visible image of sample with four pellets containing different chemicals: (1) lactose, (2) aspirin, (3) sucrose, and (4) tartaric acid.



Chemical recognition



Security



TeraView

Quality control

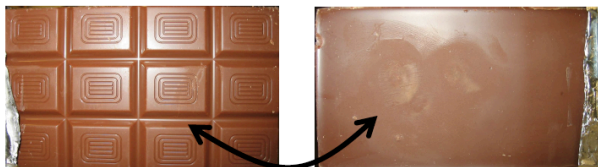
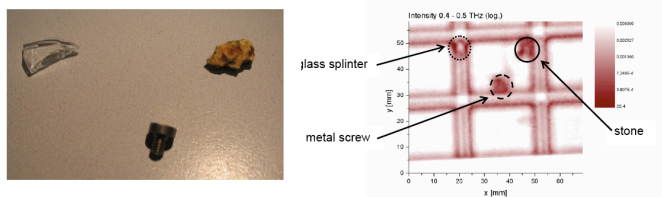
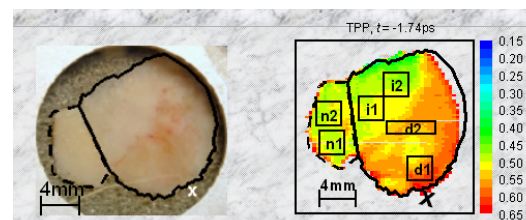


Figure 3. Front and back side of a chocolate bar after artificial contamination with a stone, a M2 metal screw and a glass splinter.



Medical imaging

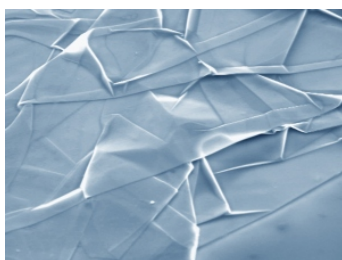
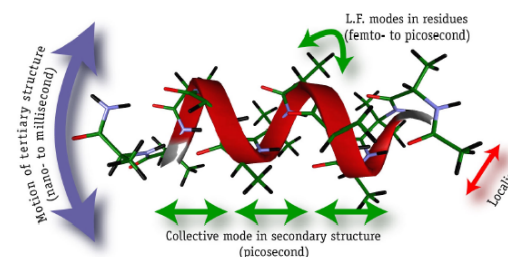




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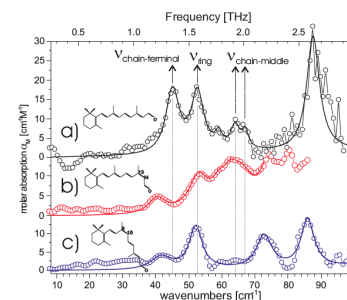
THz Spectroscopy

Superconductivity
Collective excitations
Multiferroics
Heterostructures
Metamaterials
Plasmonics

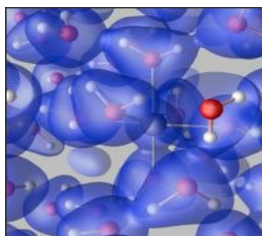


Graphene
2D chalcogenides
Black Phosphorus

Protein Folding
Amyloid fibrils
Isomers



Polar liquids
Hydrogen bonds
Van der Waals interactions
Solutions





THz sources for spectroscopy and imaging

Monochromatic Sources

Backward-Wave-Oscillators

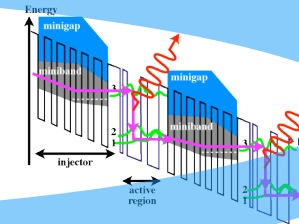


Gunn/Impatt-Diodes



Gas Lasers (CO_2 and CO_2 -pumped)
Si/Ge Lasers

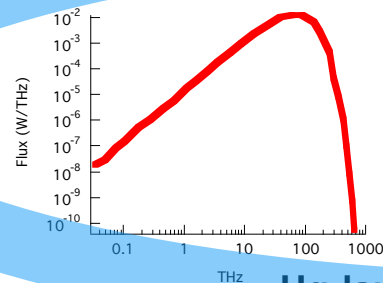
Quantum Cascade Lasers



Broadband Sources

Globar (blackbody source)

silicon carbide rod electrically heated up to 1000 to 1650°C

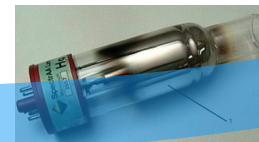


Total Flux:
10 μW between 0-1 THz
5 mW between 0-10 THz
1,5 W between 0-100 THz

Because of poor collimation
typically one has nW power at
sample at THz frequencies

Hg-lamp

Performs slightly better than Globar, below 5 THz



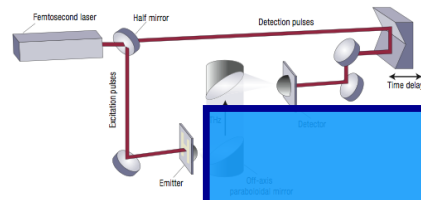


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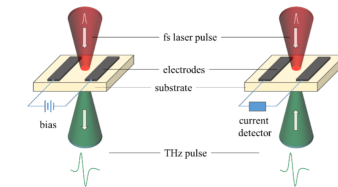
Femtosecond THz sources

0.1 THz

THz
Time Domain Spectroscopy

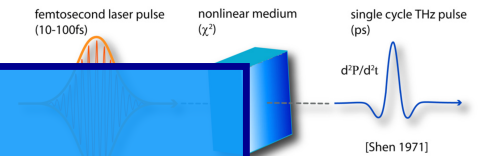


Photoconductive Antennas
GaAs, TiO₂, ...



3 THz

Optical Rectification
ZnTe, GaP, LiNbO₃, etc.
 $\eta < 10^{-2}$



Increasing THz power
---> *Increasing Laser power*
---> *Dealing with radiation damage of the emitter*
(losses in brightness/time-structure)

15 THz

Optical Parametric Amplifiers
Tunable, Narrow-Band



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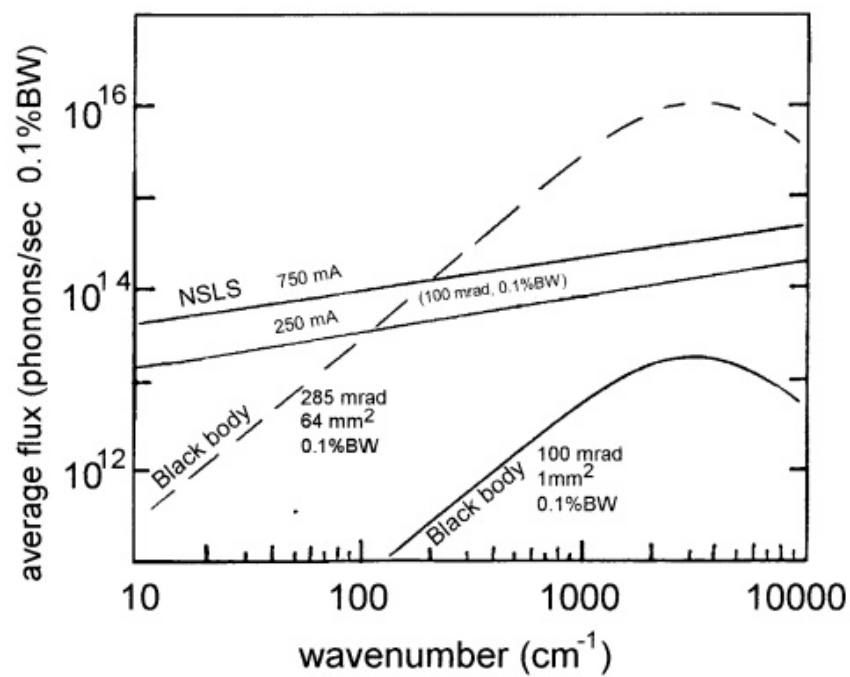
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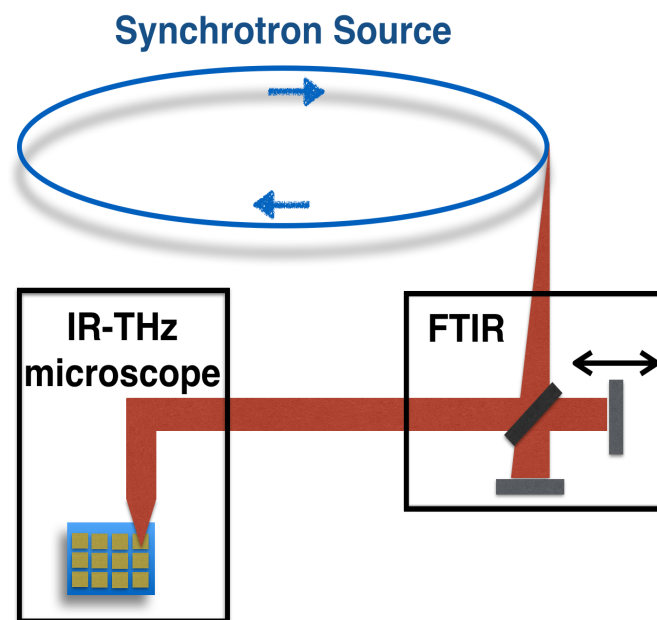


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Storage-Rings



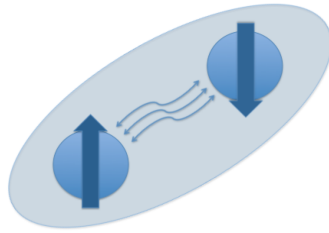
Kircher et al., JOSA B 1997



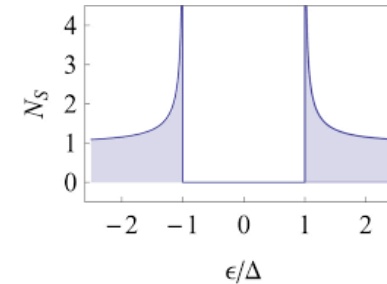


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Superconducting Gap opening



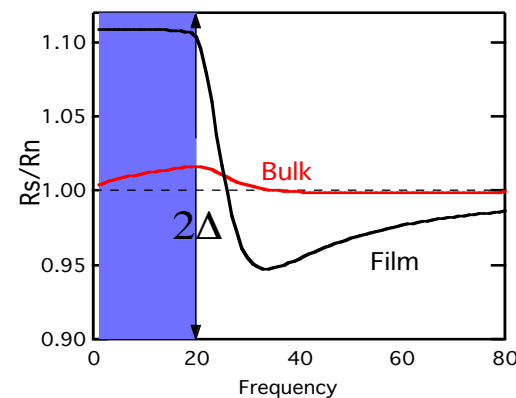
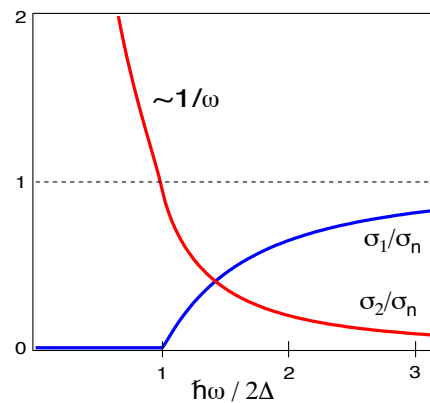
*Cooper pair formation
opens up a gap in the
density of states*



Superconductivity is ruled by *low-energy* electrodynamics:
**The Superconducting Gap size and shape provide information
on the nature and symmetry of pairing**

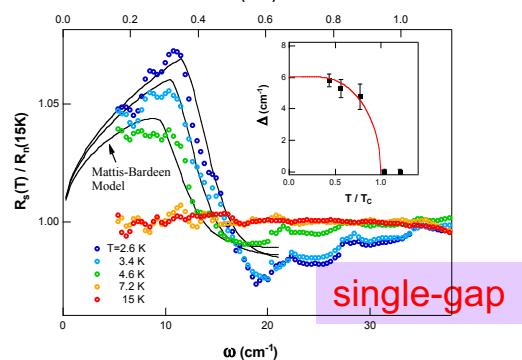
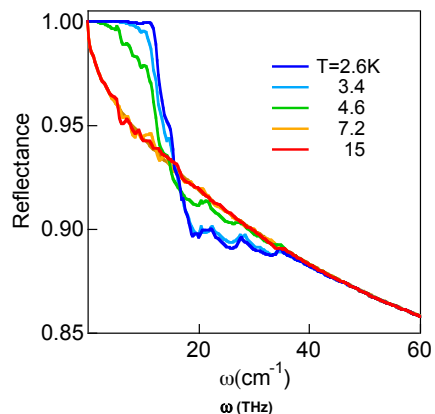
$$2\Delta / k_B T_C = 3.52 \quad \longrightarrow \quad T_c \approx 10 \text{ K} \rightarrow 2\Delta \approx 1 \text{ THz}$$

*The Mattis-Bardeen (MB) relations are derived within the
BCS theory, for a s-wave SC in the dirty limit*



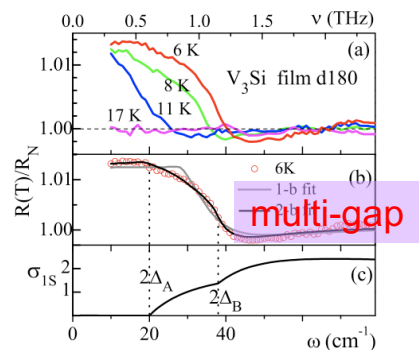
Single vs multi-gap superconductors

B-doped diamond



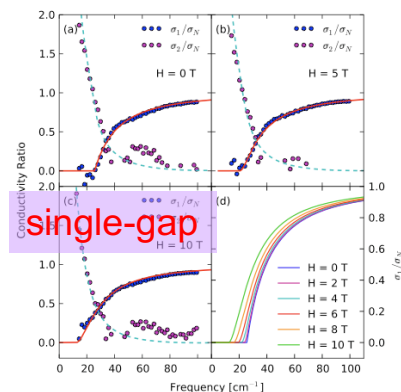
Ortolani et al., PRL 2006

V₃Si



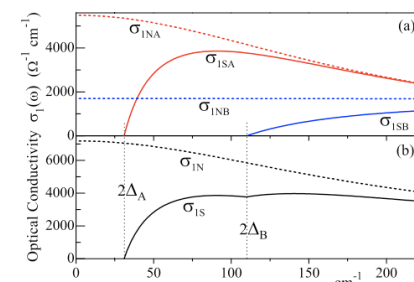
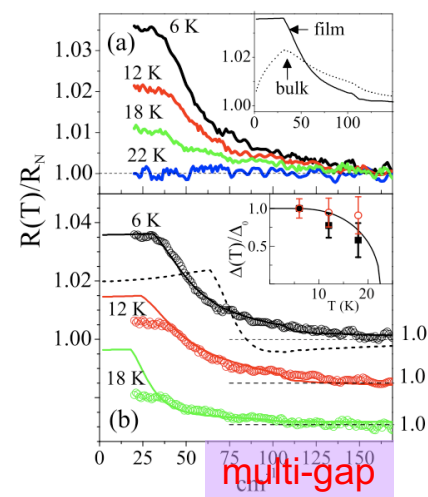
Perucchi et al., PRB 2010

Nb_{0.5}Ti_{0.5}N



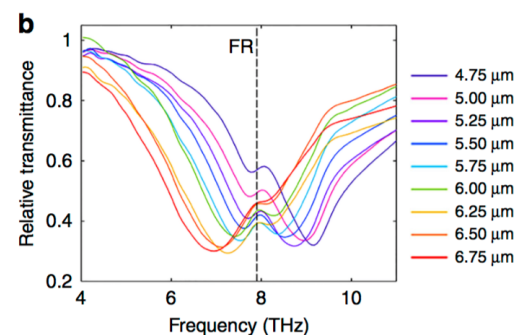
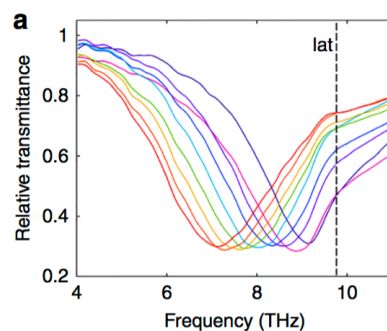
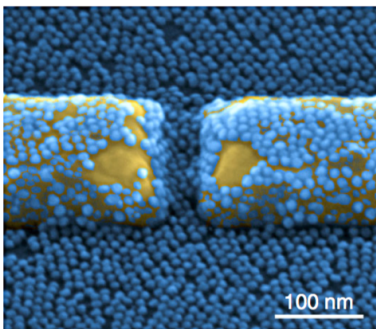
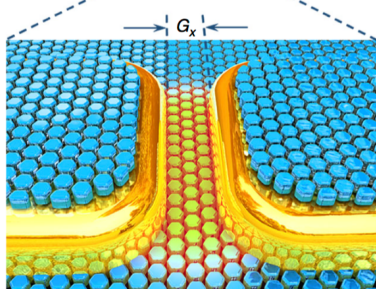
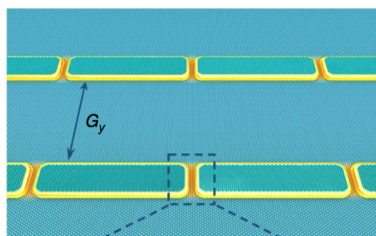
Xi et al., PRL 2011

Ba(Fe,Co)₂As₂



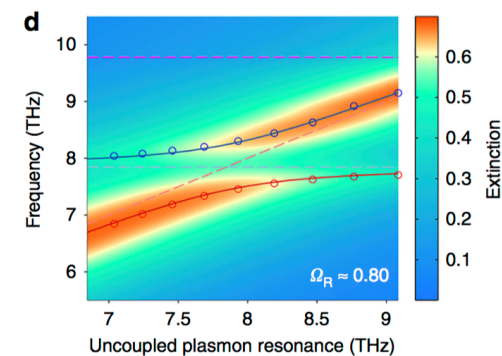
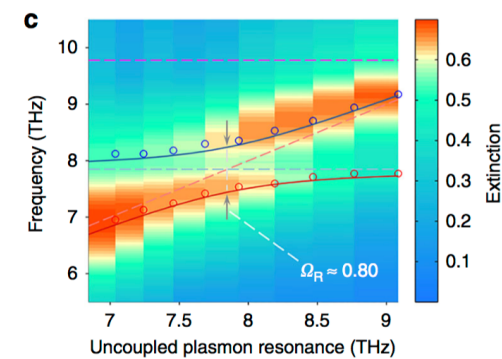
Perucchi et al., EPJB 2010

THz plasmonics



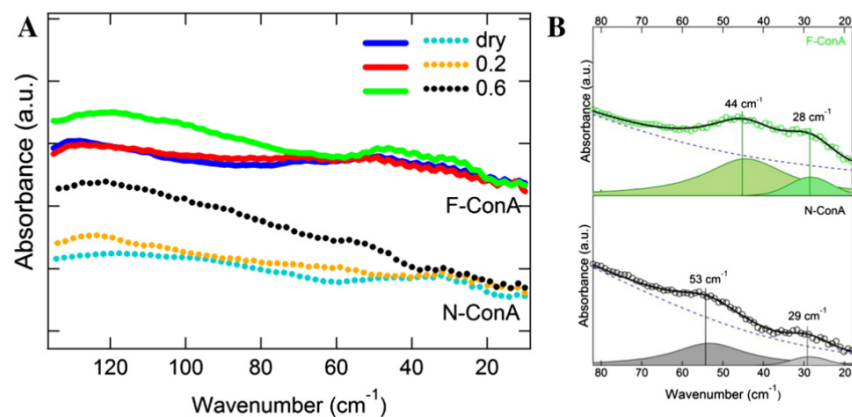
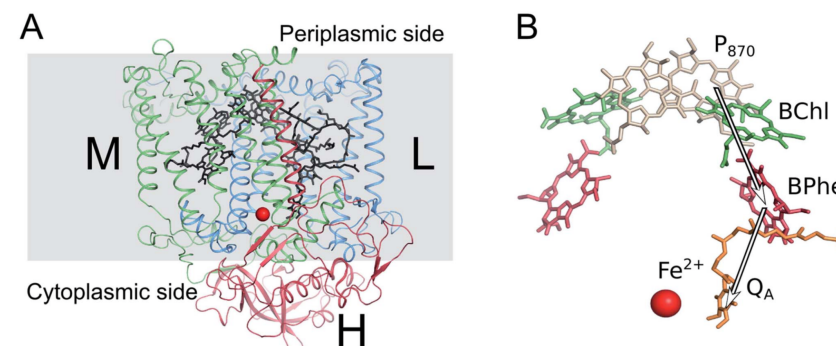
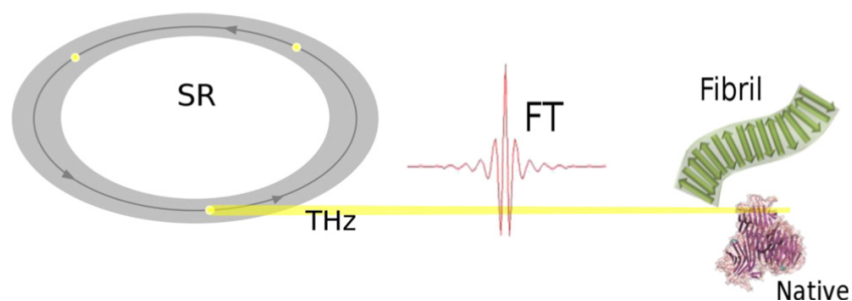
Sub-wavelength confinement (100 nm vs 50 μm)
Electric field enhancement
Increased detection

Cavity quantum electrodynamics:
Hybridization, Rabi splitting

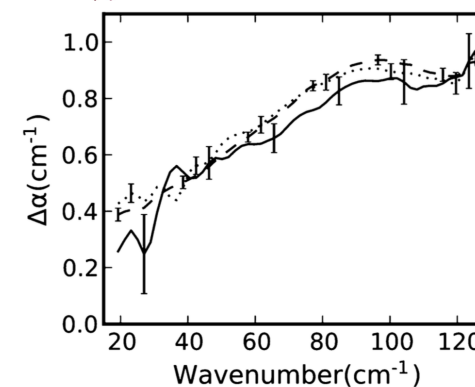


Jin et al., Nature Communications **2018**

THz spectroscopy of proteins



F. Piccirilli et al., Biophysical Chemistry **2015**



I. Lundholm et al., RSC Adv. **2014**



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THz Free Electron Lasers

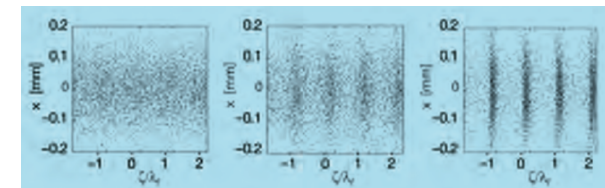
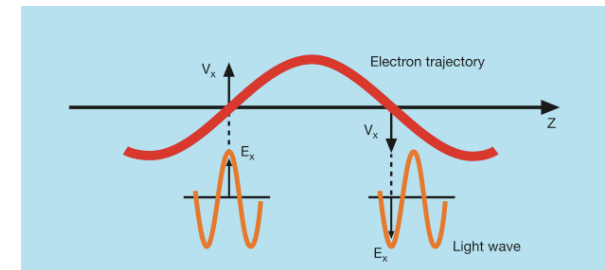
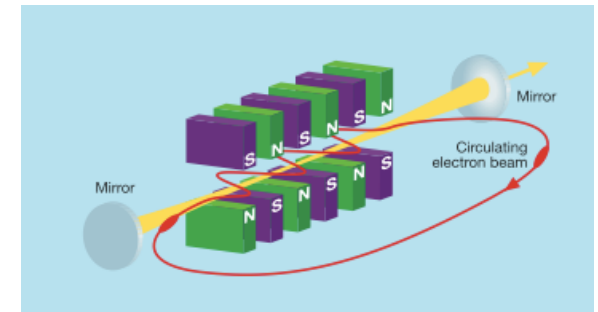
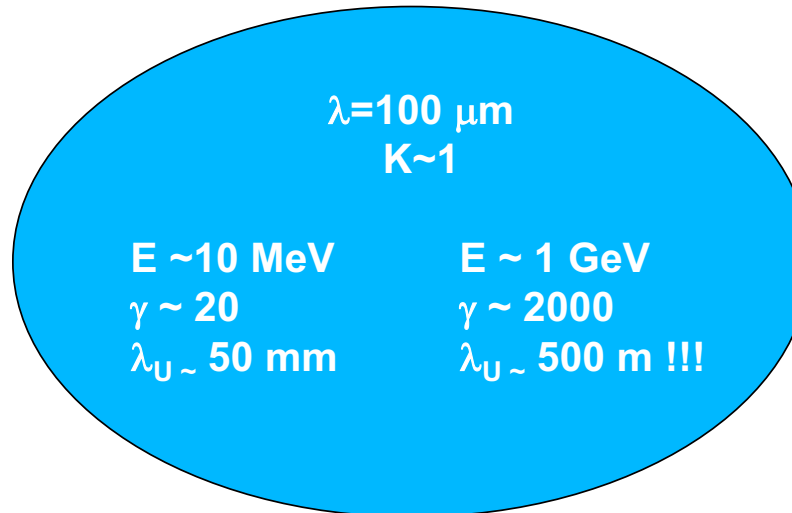
Low-gain FELs with optical cavity

- UCSB (USA)
- Budker Institute (Russia)
- CLIO (France)
- FELIX / FLARE (Netherlands)
- FELBE (Germany)

...

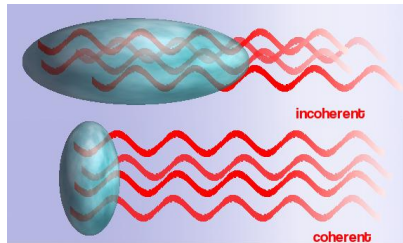
$$\lambda = \frac{\lambda_u}{2\gamma^2} \left(1 + \frac{K^2}{2} \right)$$

Upon one passage in the undulator,
radiation grows by a few percent
→ Several passages are needed
before reaching saturation



From DESY-FLASH brochure

Superradiant emission

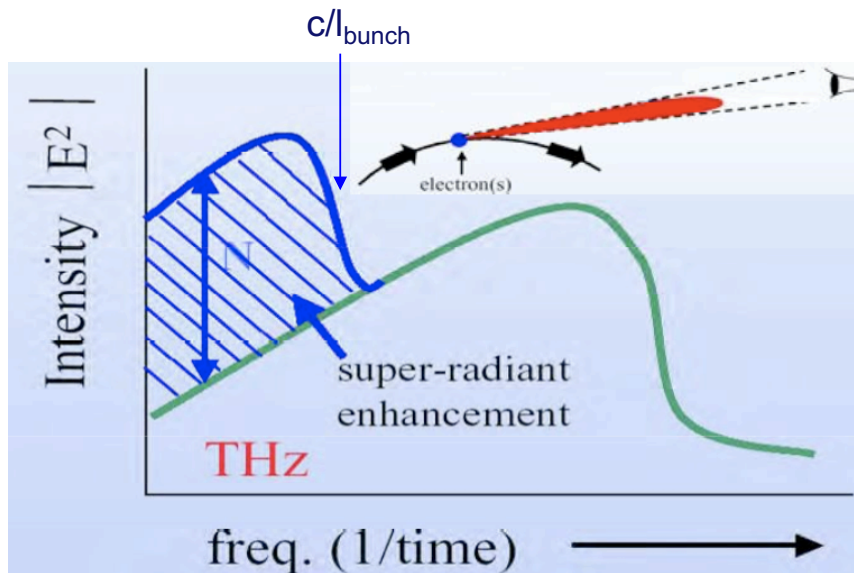


Coherent enhancement factor

$$N[1 + Nf(\omega)]$$

Form factor

$$f(\omega) = \int_{-\infty}^{+\infty} \rho(t) \exp(-i\omega t) dt$$



$N \sim 6.24 \cdot 10^7$ @ 1pC
Storage-Rings in low- α
Bessy II, Soleil, Diamond, Anka,
CLS, ...
Elettra (*bursting mode*)

$N \sim 6.24 \cdot 10^{10}$ @ 1nC
Single-pass accelerators
TeraFERMI, TELBE, FLASH

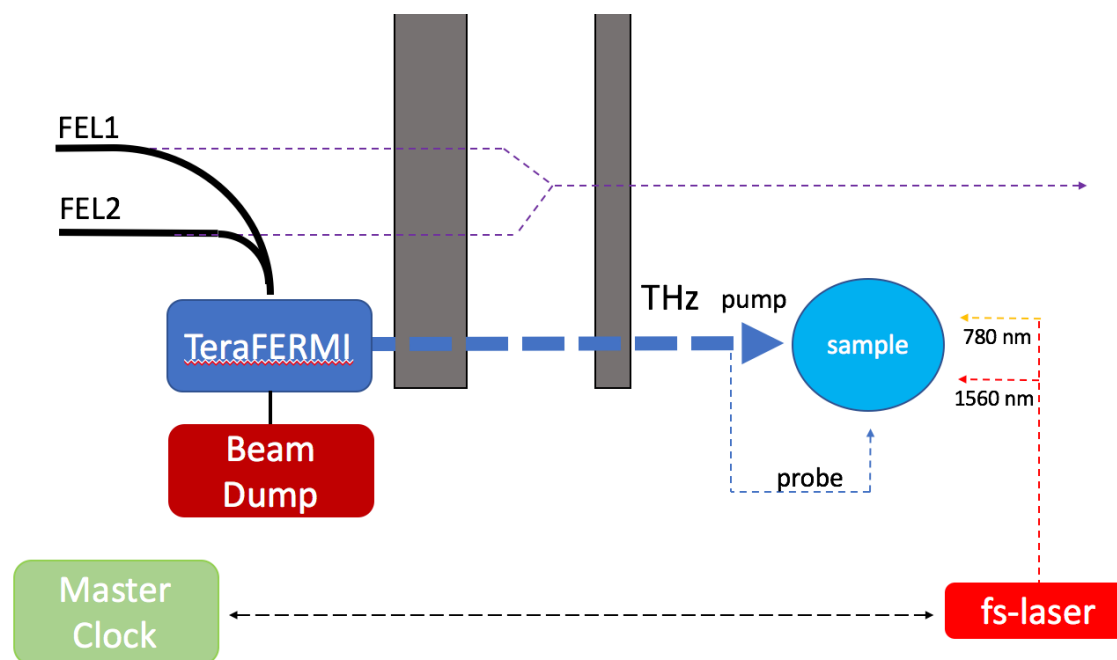
Phase Stable



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The TeraFERMI project

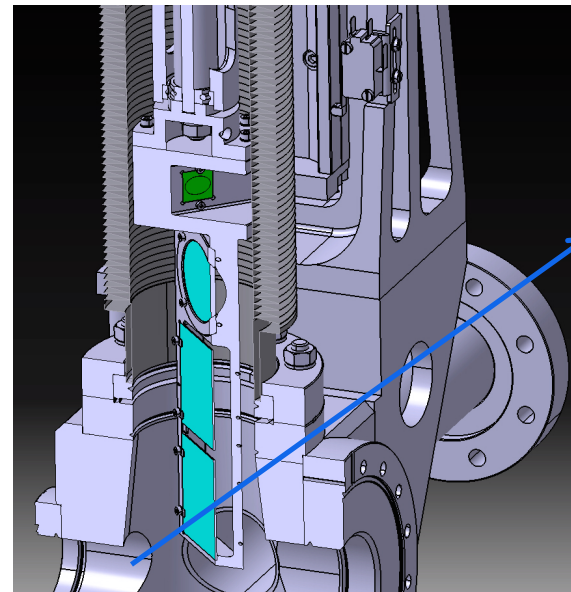
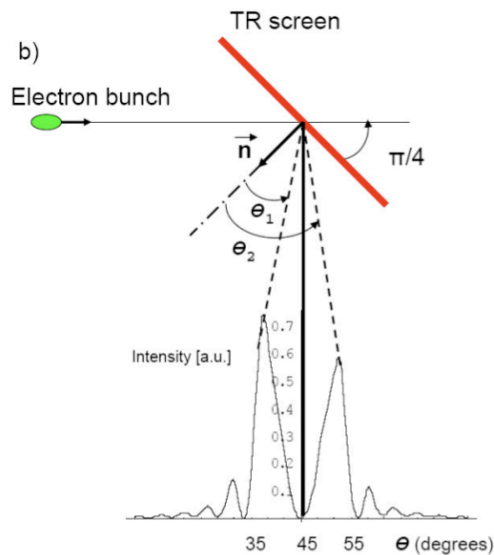
Exploiting the properties of the FERMI-FEL electron beam to produce
Short (sub-ps), Powerful ($>MV/cm$), Broadband (0.1-10 THz)
THz pulses to be used as a **Pump** beam for ultrafast nonlinear spectroscopies





Coherent Transition Radiation

Transition Radiation occurs when relativistic electrons cross the boundary between two media of different dielectric constant



The Ginzburg-Frank equation:

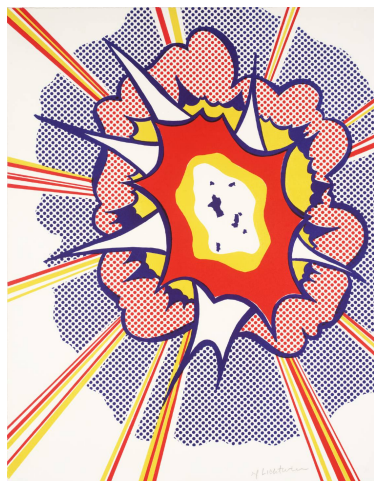
$$\frac{d^2U}{d\omega d\Omega} = \frac{e^2}{4\pi^3 \epsilon_0 c} \frac{\beta^2 \sin^2 \theta}{(1 - \beta^2 \cos^2 \theta)^2}$$



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Why THz?

THz light couples to electronic, vibrational and magnetic excitations



Optical pump



THz pump

THz control of matter

Nonlinear Phononics

Superconductivity
Metal-insulator
transitions
Ferroelectrics

Macromolecules

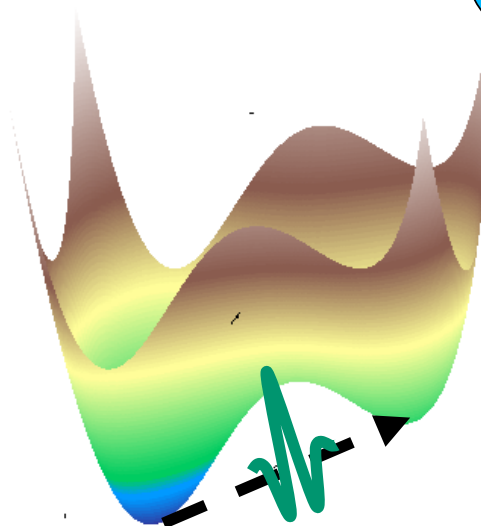
Protein folding
RNA/DNA

Molecular alignment

Polar liquids

Reaction pathways

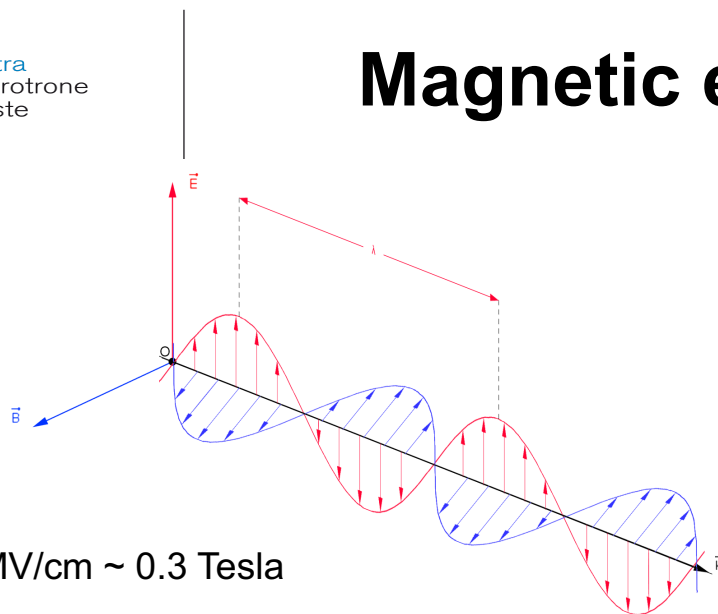
catalysis



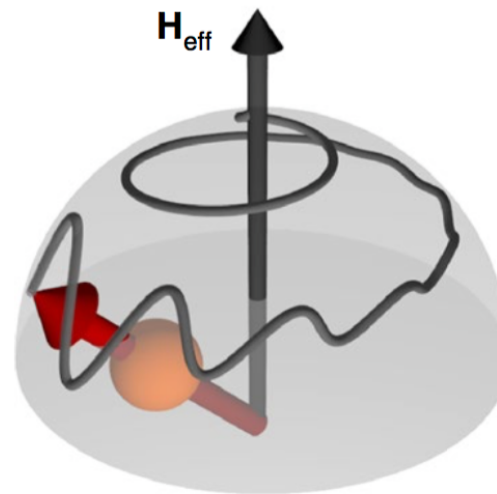


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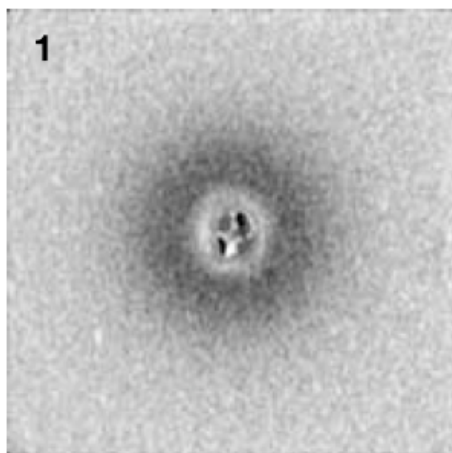
Magnetic excitations



1 MV/cm ~ 0.3 Tesla



Neeraj, 2021



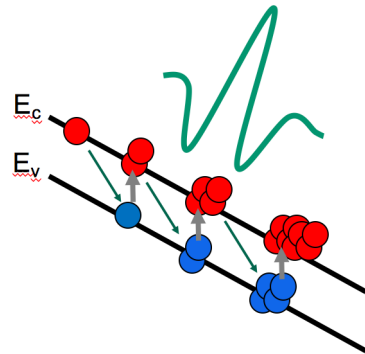
Tudosa, 2004

Ultrafast magnetic
switching

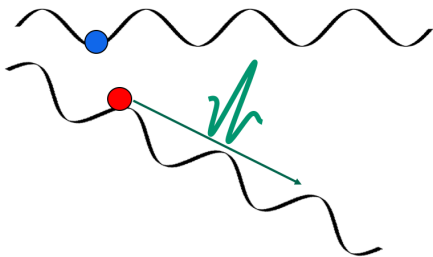


Electronic excitations

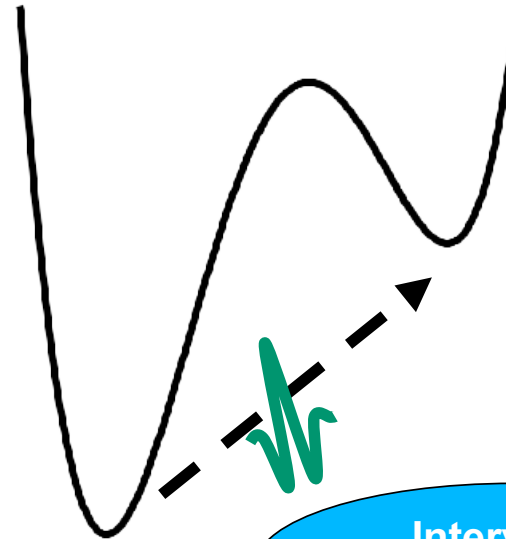
Impact Ionization



Zener tunneling



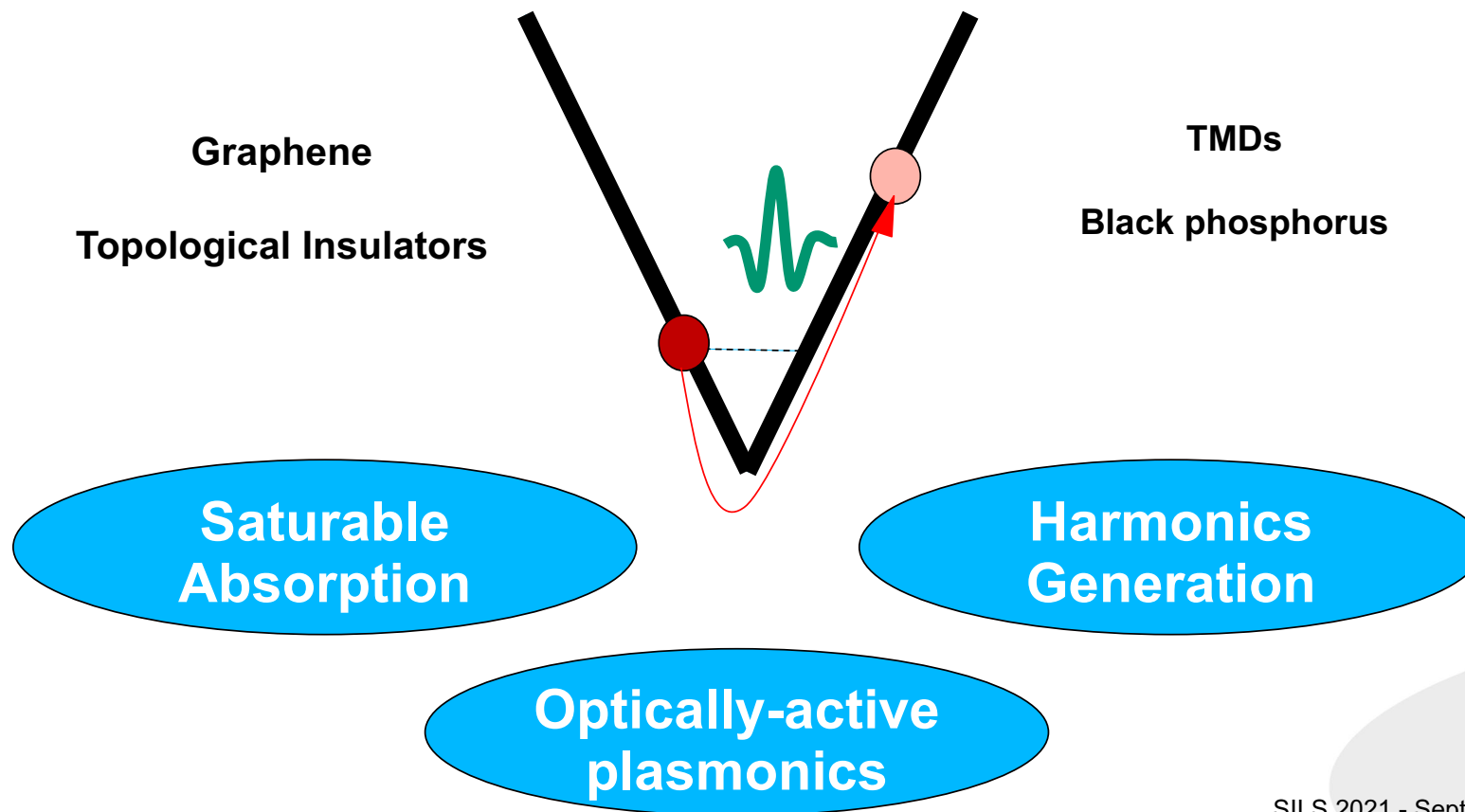
Intervalley scattering



Franz-Keldish



Dirac materials





Nonlinear susceptibility

THz pump can resonantly excite electrons or lattice modes

linear

Pockels

Kerr

$$P_i = \varepsilon_0(\chi_{ij}^{(1)}E_j + \chi_{ijk}^{(2)}E_jE_k + \chi_{ijkl}^{(3)}E_jE_kE_l + \dots)$$



THz pump-IR probe

THz pump can resonantly excite electrons or lattice modes

linear

Pockels

Kerr

$$P_i = \varepsilon_0 (\chi_{ij}^{(1)} E_j + \chi_{ijk}^{(2)} E_j E_k + \chi_{ijkl}^{(3)} E_j E_k^2 + \dots)$$

E_j = Probe (780 nm)

E_k = Pump (THz)

Pockels effect:

An external electric field induces linear variation of the refractive index on a material
Centro-symmetric crystals can not exhibit this effect (no preferred direction)

Kerr effect:

The change of the refractive index is proportional to the square of the applied electric field.



THz fluence-dependent spectroscopy

THz pump can resonantly excite electrons or lattice modes

linear

Pockels

Kerr

$$P_i = \varepsilon_0 (\chi_{ij}^{(1)} E_j + \chi_{ijk}^{(2)} E_j^2 + \chi_{ijkl}^{(3)} E_j^3 + \dots)$$

E_j = THz field

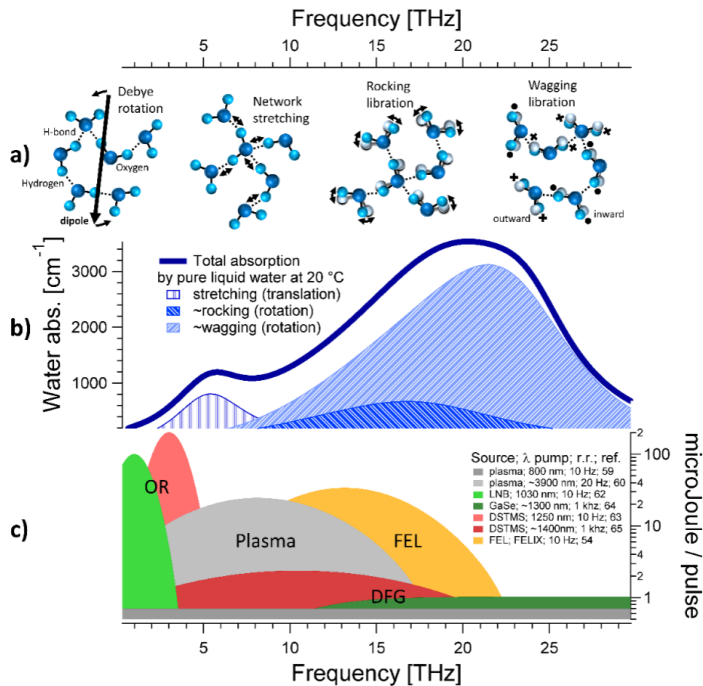
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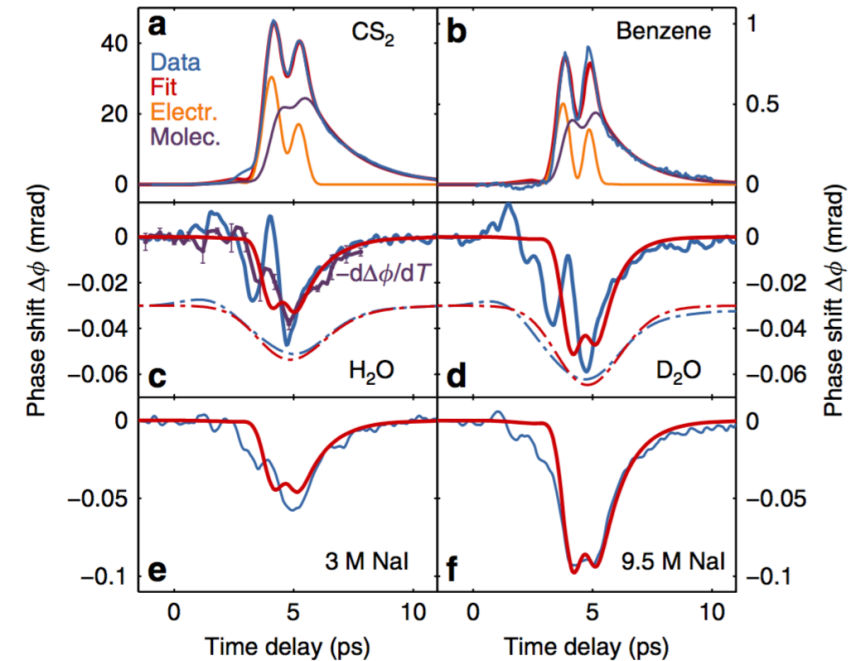
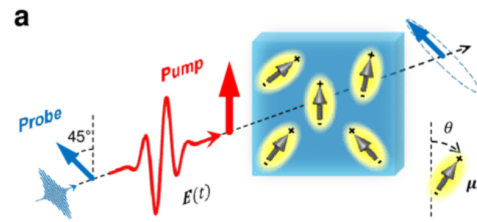
Kerr effect:

The change of the refractive index is proportional to the square of the applied electric field.

THz Kerr effect in water



F. Novelli et al., Materials 2020

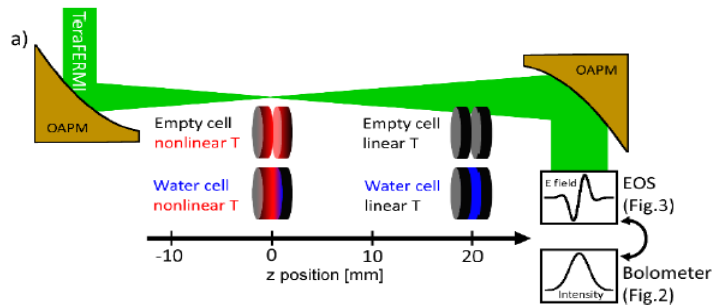


P. Zalden et al., Nat. Comm. 2018

Single-cycle electromagnetic pulses in the THz regime orient the dipole moments of liquid water along their electric field. Given the resulting negative sign of $\Delta n = n_{\parallel} - n_{\perp} < 0$, the polarizability of water molecules in the liquid state is lower parallel to their dipole moment than perpendicular, i.e., $\alpha_{\parallel} < \alpha_{\perp}$

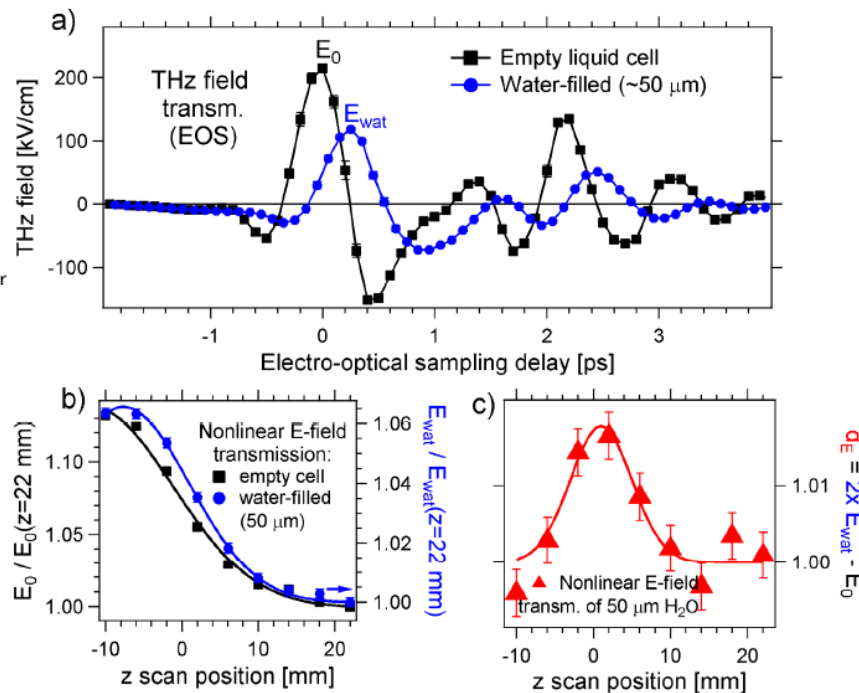
THz-fluence dependent measurements

Nonlinear THz transmission by liquid water



By open aperture z-scan transmission experiments on static liquid cells we detect the THz fields with electro-optical techniques.

We show that it is possible to quantify the **nonlinear response of liquid water at ~1 THz** even when large signals originate from the sample holder windows.



$$\alpha(P) = \alpha + \alpha^{NL} \times P,$$

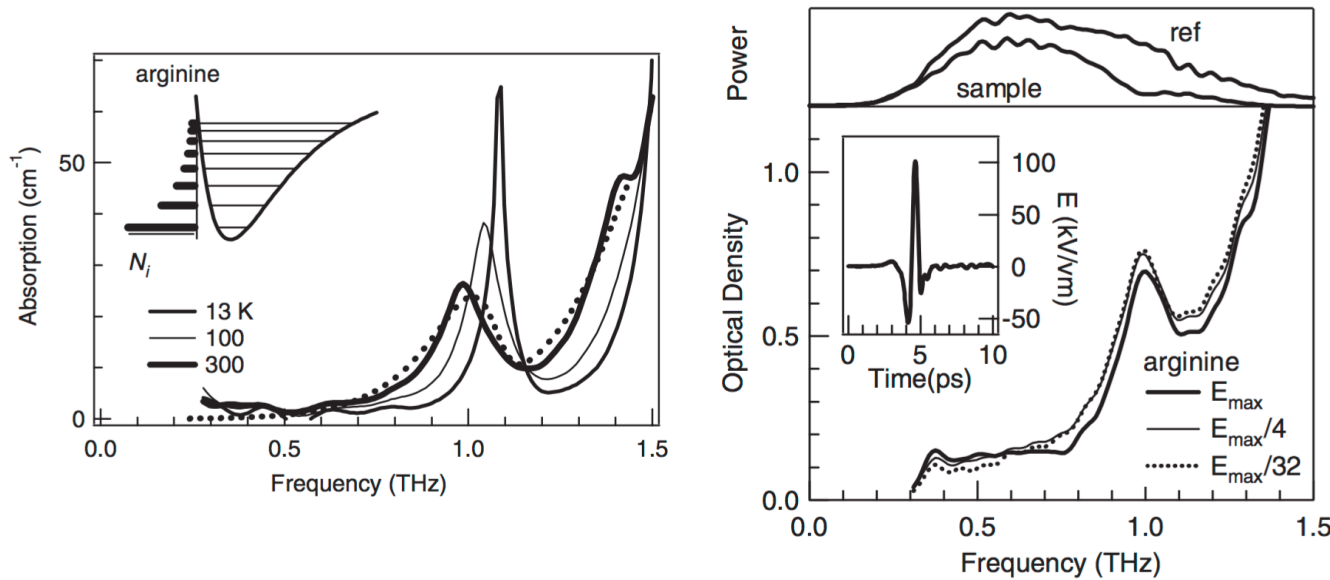
$$\alpha^{NL} \approx (1 - d_I(0)) \frac{\alpha}{P \times (1 - e^{-\alpha L})},$$

$$n(P) = n + n^{NL} \times P,$$

$$n^{NL} = \Delta T \frac{\alpha}{P \times (1 - e^{-\alpha L})} \frac{\lambda}{1.8 \times (1 - S)^{0.25}},$$

THz-fluence dependent measurements

Ladder climbing of the anharmonic intermolecular potential in aminoacids



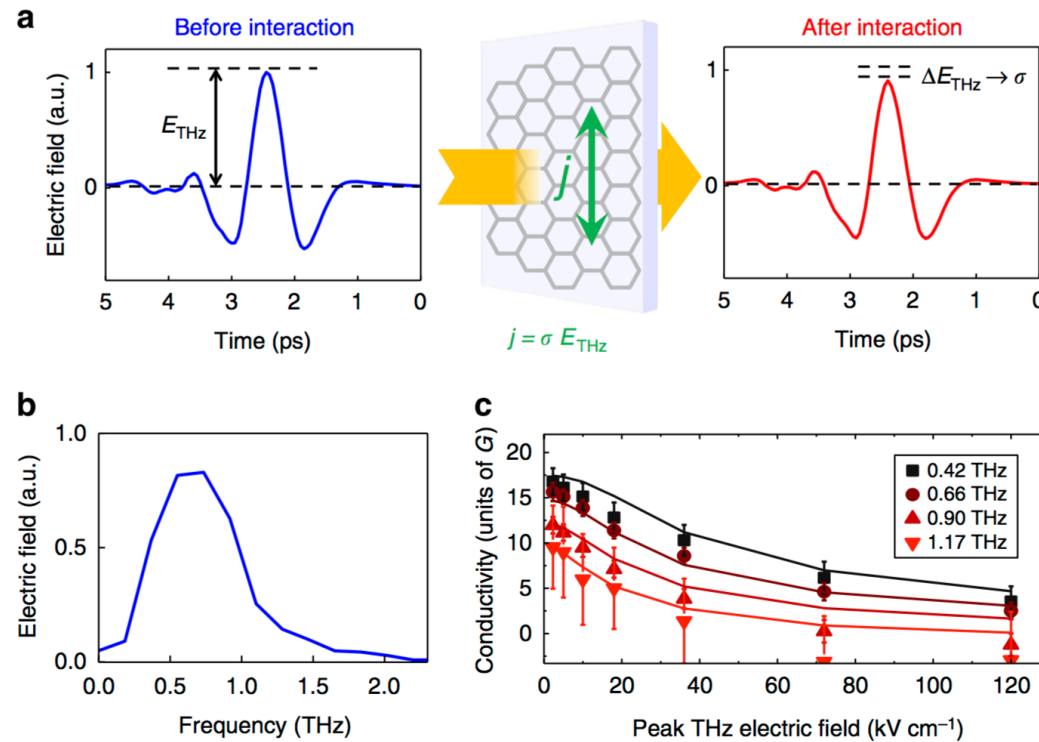
$$\begin{aligned} \frac{d\rho_{i,i}(t)}{dt} = & -\gamma_1(\rho_{i,i} - \rho_{i+1,i+1}) \\ & -i\frac{E(t)}{\hbar}(\mu_{i-1,i}\rho_{i,i-1} - \mu_{i-1,i}^*\rho_{i,i-1}^*) \\ & +i\frac{E(t)}{\hbar}(\mu_{i,i+1}\rho_{i+1,i} - \mu_{i,i+1}^*\rho_{i+1,i}^*), \end{aligned} \quad (1)$$

$$\begin{aligned} \frac{d\rho_{i+1,i}(t)}{dt} = & \frac{d\rho_{i,i+1}^*(t)}{dt} \\ = & -\left(i2\pi\nu_{i,i+1} + i(\mu_{i+1,i+1} - \mu_{i,i})\frac{E(t)}{\hbar} + \gamma_2\right) \\ & \times \rho_{i+1,i} - i\frac{E(t)}{\hbar}\mu_{i+1,i}(\rho_{i+1,i+1} - \rho_{i,i}), \end{aligned} \quad (2)$$

“The mode population around the levels with $\langle n \rangle = 2.2$ rises to a higher level with $\langle n \rangle = 22.2$, indicating ladder climbing of 20 steps. This is far from thermal equilibrium.”



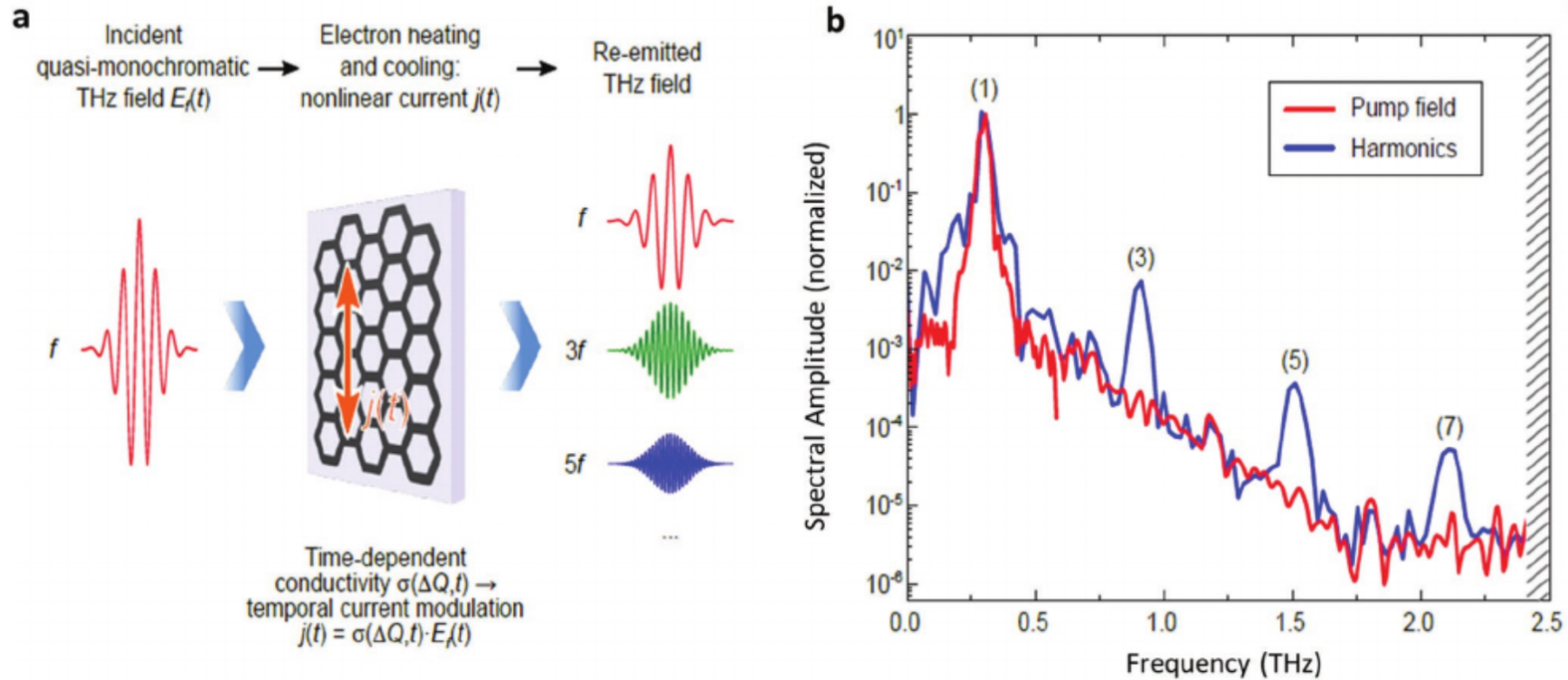
THz field induced transparency in graphene



Mics et al., Nature Communication (2015)



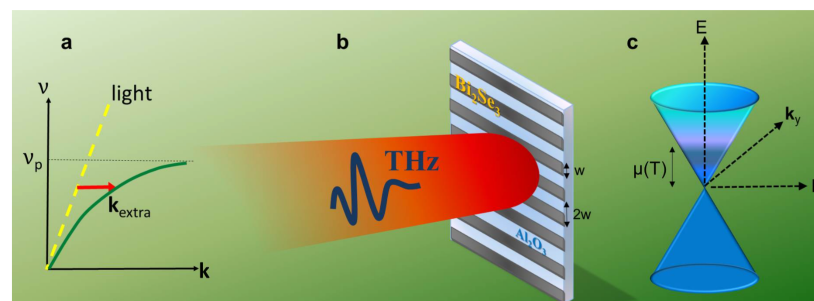
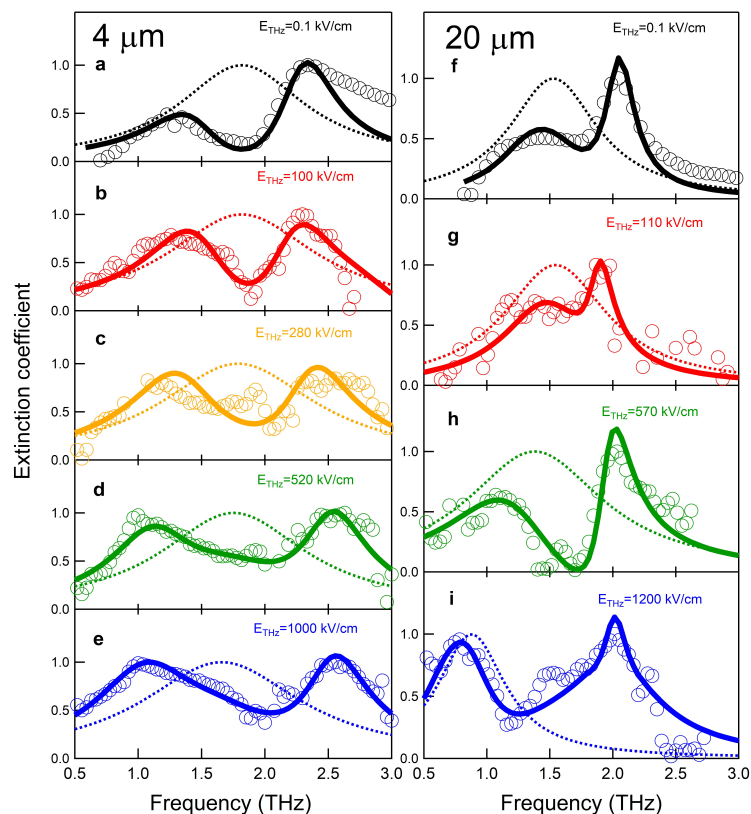
THz Harmonics generation in graphene



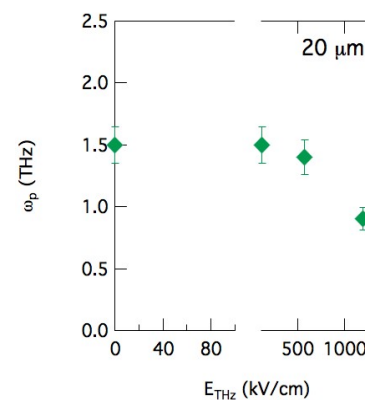
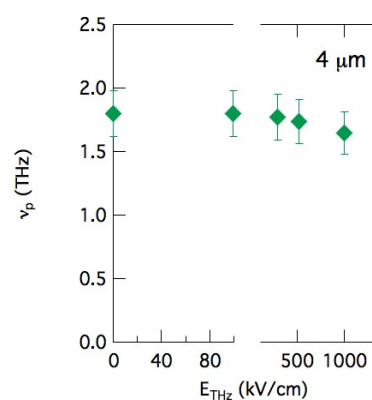
Hafez et al., Nature (2018)

THz-fluence dependent measurements

Tuning plasmons in Bi2Se3



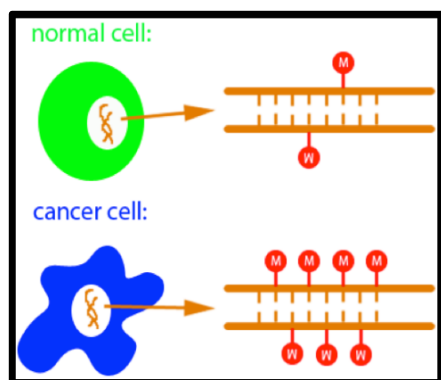
$$\nu_p = \sqrt{\frac{\pi}{w}} \left(\frac{e^2}{4\pi\epsilon_0\epsilon\hbar} v_F \sqrt{2\pi n g_s g_v} \right)^{1/2},$$



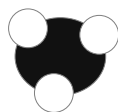
P. Di Pietro et al., Phys. Rev. Lett. (2020)

THz irradiation

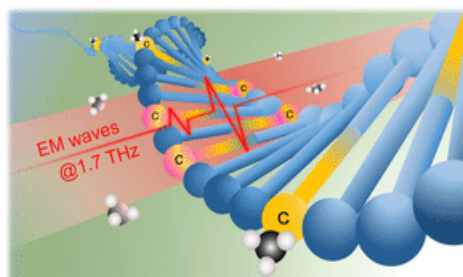
THz induced demethylation



Methyl groups



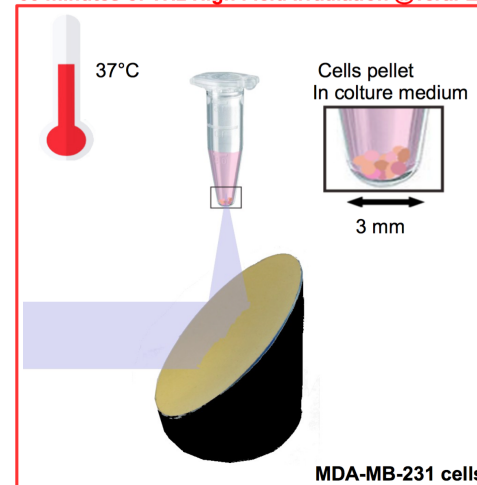
Cancer cells can show enhanced DNA Methylation



Intense resonant THz radiation induces DNA demethylation

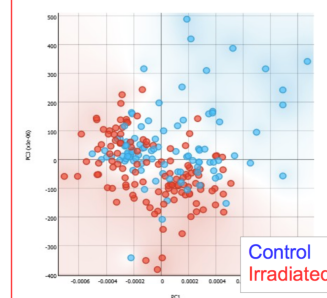
Treatment

60 minutes of THz High Field Irradiation @TeraFERMI



Probe

SR-IR @SISSI on single cells



F. Piccirilli, M. Pachetti et al. (2021)

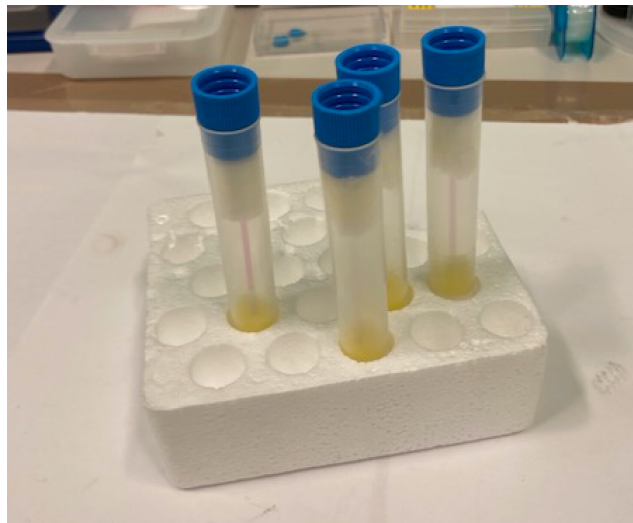
THz irradiation

Covid inactivation through THz irradiation

Elettra - Università di Venezia Cà Foscari - ICGEB

RNA modes resonate at THz frequencies → Inducing RNA damage while keeping the capsid intact → vaccines

Most plastics are transparent to THz radiation.
Unbreakable sample holders can be used.
Virus is sealed within 2 plastic containers.



Experiments are still ongoing...



Conclusions

THz gap is nowadays filled and THz spectroscopy can be performed with several different conventional equipments

Storage Rings provide bright sources of THz radiation for steady state studies

Small samples, Microscopy, Extreme conditions, etc.

THz control of matter requires short/powerful pulses
(*sub-ps duration / >MW peak power / ~MV/cm electric fields*)

FEL-based THz beamlines can produce THz pulses for non-linear studies in the whole THz range



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THANK YOU

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