



LEDS2023

Longitudinal Electron beam Dynamics for coherent light Sources



Compact



Beam Loading Effects in photo-injectors. Simulations and Measurements

Longitudinal Electron beam Dynamics for coherent light Sources

Javier Olivares Herrador (javier.olivares.herrador@cern.ch), Andrea Latina

03.10.2023



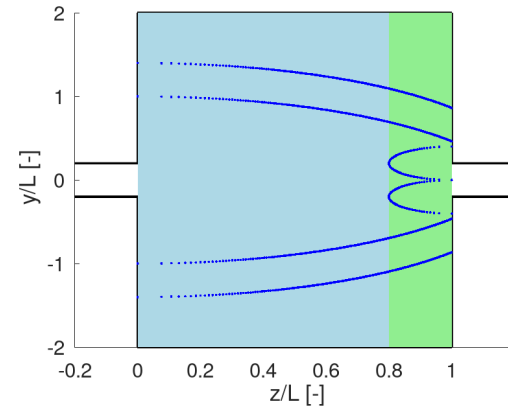
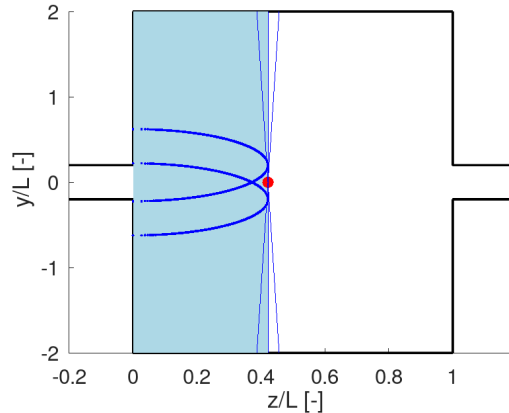
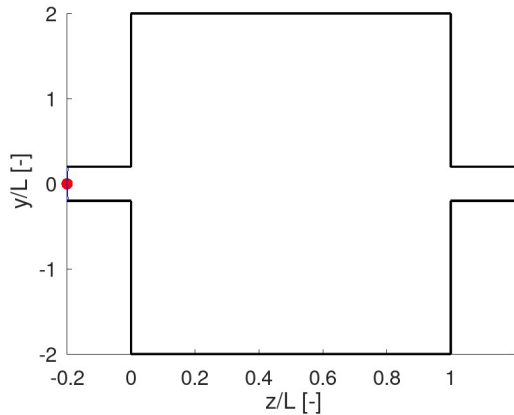
Outline

- **PART I:** Introduction. The Beam Loading Effect
 - TW model
- **PART II:** BL in Standing-Wave Photo-injectors. Particularities
- **PART III:** BL Simulations with RF-Track
- **PART IV:** Results
 - Simulations vs Measurements in the CLEAR facility
 - BL Compensation

PART I: The Beam Loading Effect

Beam Loading Effect

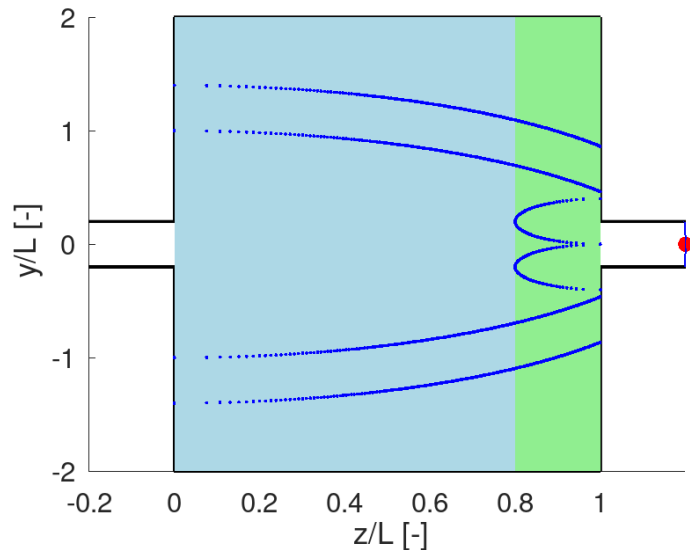
- Beam – Cavity interaction: **Excitation** of the **fundamental accelerating mode**
 - Beam-induced field causes **deceleration**



> Animation of a particle flying through a cavity and leaving an EM field behind it.

Beam Loading Effect

- Beam – Cavity interaction: **Excitation** of the **fundamental accelerating mode**
 - Beam-induced field causes **deceleration**

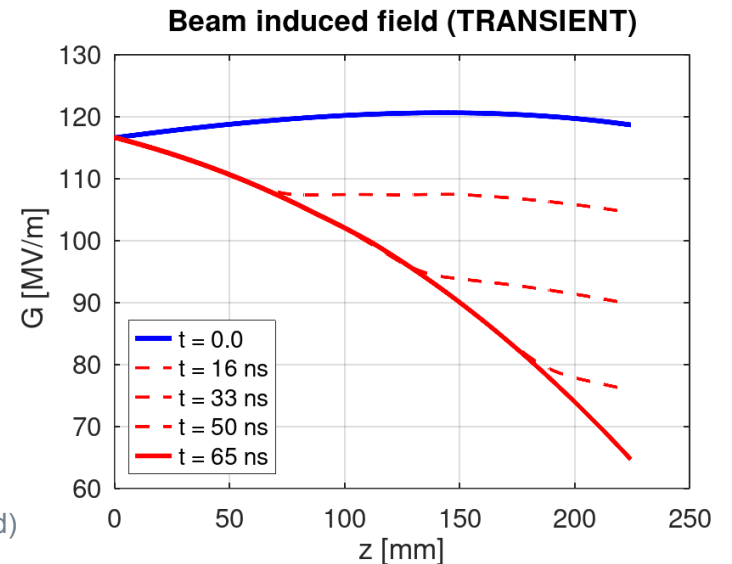


- The induced excitation lasts for a long time
 - **Long range effect**
 - Accumulated from bunch-to-bunch

> Animation of a particle flying through a cavity and leaving an EM field behind it.

Beam Loading Effect

- **What:** Reduction of available accelerating gradient
- **Origin:** Beam – Cavity interaction
- **Consequences:** Transient response
 - Different energy loss from bunch to bunch
- **Motivation:** High I, Compact accelerating structures



[1] P. Lapostolle. *Linear Accelerators*. North Holland Publishing Company, 1970 (Amsterdam, Holland)

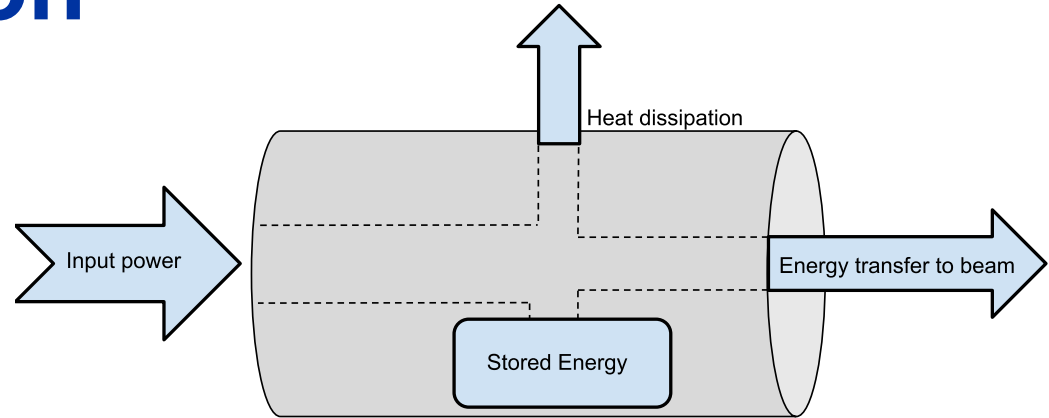
[2] A. Grudiev, A. Lunin, V. Yakovlev. *Analytical solutions for transient and steady state beam loading in arbitrary travelling wave accelerating structures*. Phys. Rev. Special topics **14**, 052001 (2011)

> Theoretical analysis of beam loading effect based on CLIC's main linac [2]

I. Energy Conservation

- Poynting Theorem

$$-\underbrace{\frac{\partial u(\vec{r}, t)}{\partial t}}_{\text{Stored EM energy density variation}} = \underbrace{\vec{\nabla} \cdot \vec{S}(\vec{r}, t)}_{\text{Power Flow \& Loss}} + \underbrace{\vec{E}(\vec{r}, t) \cdot \vec{J}(\vec{r}, t)}_{\text{Field-Beam Interaction}}$$



> Energy balance schematics for an accelerating structure

- Figures of merit:

- Group velocity

$$v_g = \frac{P_{\text{flow}}}{w} [\text{m/s}]$$

- Quality factor

$$Q = \omega_{\text{RF}} \frac{w}{p_{\text{diss}}}$$

- Shunt impedance (p.u.l)

$$r_e = \frac{G_{\text{eff}}^2}{p_{\text{diss}}} [\Omega/\text{m}]$$

[3] Thomas P. Wangler. *RF linear accelerators*. Wiley-VCH 2008 (Amsterdam, Holland)

I. Gradient description

- **Gradient:** Averaged E-field *affecting* the particle $t_f(z) \simeq \frac{z}{c}$

TW wave $E_z(z, t) = \text{Re}[E_0(z, t)e^{j(kz-\omega t)}]$ leads to $G(z, t)$

SW wave $E_z(z, t) = \text{Re}[E_0(z, t)e^{j\omega t}]$ leads to $G_{\text{eff}}(z, t)$

- Time Transit factor

$$\mathcal{T}(z, t) = \frac{G_{\text{eff}}(z, t)}{G(z, t)} = \frac{\frac{1}{L} \int_0^L \text{Re}[E_0(z, t)e^{j\omega t(z)}] dz}{\frac{1}{L} \int_0^L |E_z(z, t)| dz}$$

[3] Thomas P. Wangler. *RF linear accelerators*. Wiley-VCH 2008 (Amsterdam, Holland)

I. Power-Diffusion PDE

- From Poynting: Equation in terms of Gradient:

$$-\frac{\partial G_{\text{eff}}}{\partial t} = v_g \frac{\partial G_{\text{eff}}}{\partial z} + \left(-\frac{v_g Q}{r_{\text{eff}}} \frac{\partial(r_{\text{eff}}/Q)}{\partial z} + \frac{\omega}{Q} + \frac{\partial v_g}{\partial z} \right) \frac{G_{\text{eff}}}{2} + \underbrace{\frac{\omega r_{\text{eff}} \tilde{I}}{2Q}}_{\text{Beam Loading term!}}$$

Some features:

- **Paraxial + Quasi-static** approximation
- **Beam Loading term: Decelerating** gradient dependent on **Intensity**.
- Assumes **causality!**
- Matches [2] for the TW ultrarelativistic case

[4] J. Olivares Herrador, D. Esperante Pereira, N. Fuster, B. Gimeno, and A. Latina, “Beam Loading Simulation for Relativistic and Ultrarelativistic Beams in the Tracking Code RF-Track”, in Proc. LINAC’22, Liverpool, UK, Aug.-Sep. 2022, pp. 569–572. doi:10.18429/JACoW-LINAC2022-TUPORI13

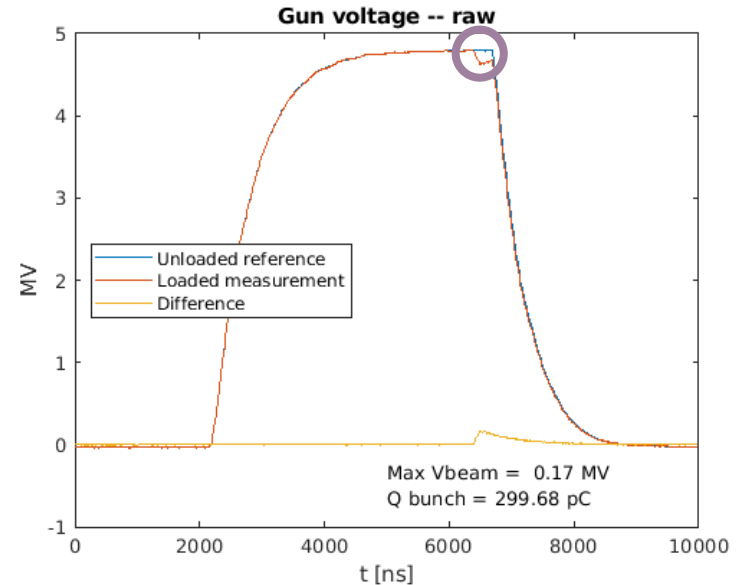
PART II: BL in SW photo-injectors

SW Beam Loading Model

- Photo-Injector = SW accelerating structure $\rightarrow v_g = 0$ m/s.
- No power flow \rightarrow Input power has to be considered

$$\frac{\partial G_{\text{eff}}}{\partial t} = -\frac{\omega}{Q} \frac{G_{\text{eff}}}{2} - \frac{\omega r_{\text{eff}} \tilde{I}}{2Q} + \frac{G_{\text{target}} \omega}{2Q}$$

Gradient reduction takes place while cavity is being fed



> Voltage along an RF-cycle for CLEAR SW photo-injector. Measurement.

SW Beam Loading Model

- Photo-Injector = SW accelerating structure $\rightarrow v_g = 0$ m/s.
- No power flow \rightarrow Input power has to be considered
- $v < c$ in the beginning ...

- Phase slippage \rightarrow Already taken into account with effective figures of merit $t_f(z) = \int_0^z \frac{dz'}{\beta(z')c}$

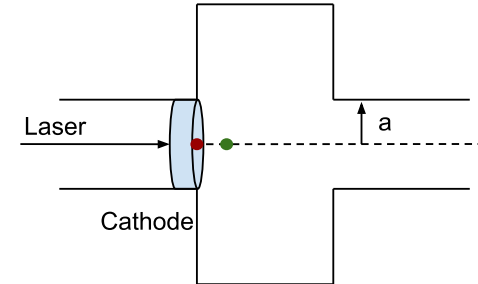
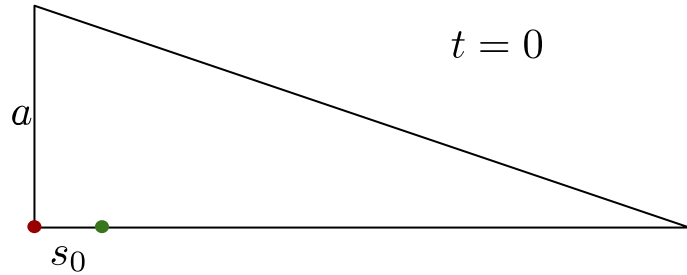
SW Beam Loading Model

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- $v < c$ in the beginning ...
 - Phase slippage \rightarrow Already taken into account with effective figures of merit $t_f(z) = \int_0^z \frac{dz'}{\beta(z')c}$
 - Causality?
 - Does the wake of the preceding particle affect the particle ahead?

To be addressed in each specific case

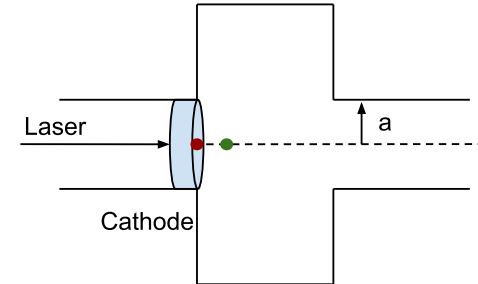
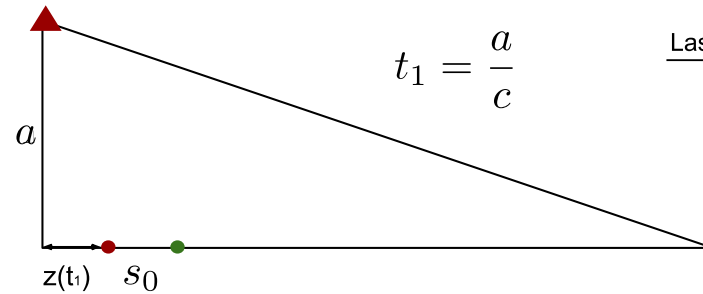
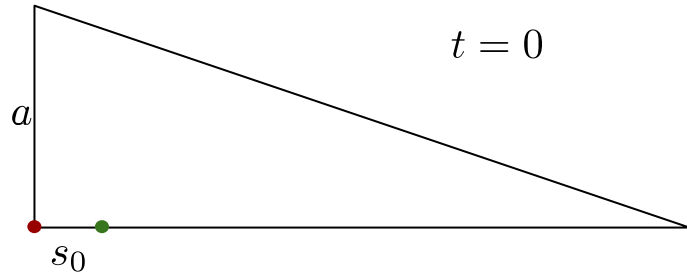
SW Beam Loading Model: Causality

- **Causality:** Does the wake of the preceding particle affect the particle ahead?



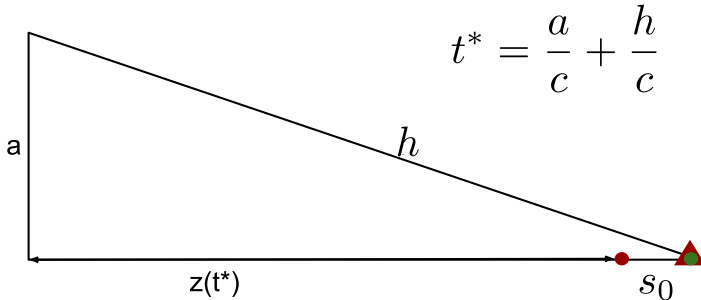
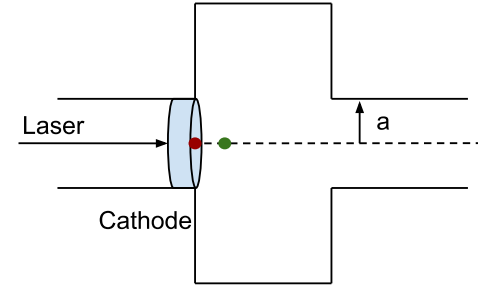
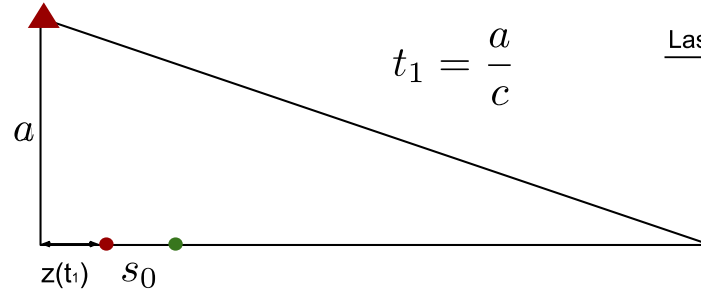
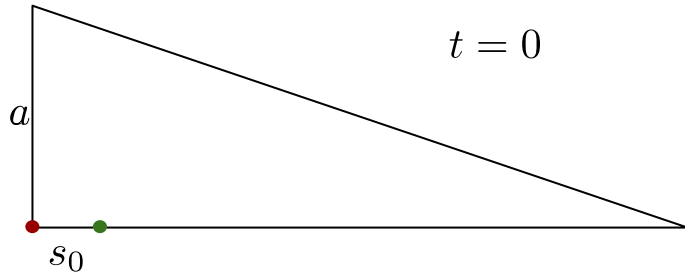
SW Beam Loading Model: Causality

- Causality:** Does the wake of the preceding particle affect the particle ahead?



SW Beam Loading Model: Causality

- Causality:** Does the wake of the preceding particle affect the particle ahead?



$$\begin{cases} z(t^*) = \int_0^{t^*} \beta(t)c dt \\ h = \sqrt{a^2 + (s_0 + z(t^*))^2} \end{cases}$$

$$\exists t^* \in \mathbb{R} \text{ so that } t^* = \frac{a}{c} + \frac{1}{c} \sqrt{a^2 + (s_0 + z(t^*))^2}$$

Catch-up condition

SW Beam Loading Model: Causality

- **Causality:** Does the wake of the preceding particle affect the particle ahead?
- CLEAR facility @ CERN

Magnitude	Units	Value
a	mm	10
$\langle G \rangle$	MV/m	30.2
σ_z	mm	0.1 – 1.2

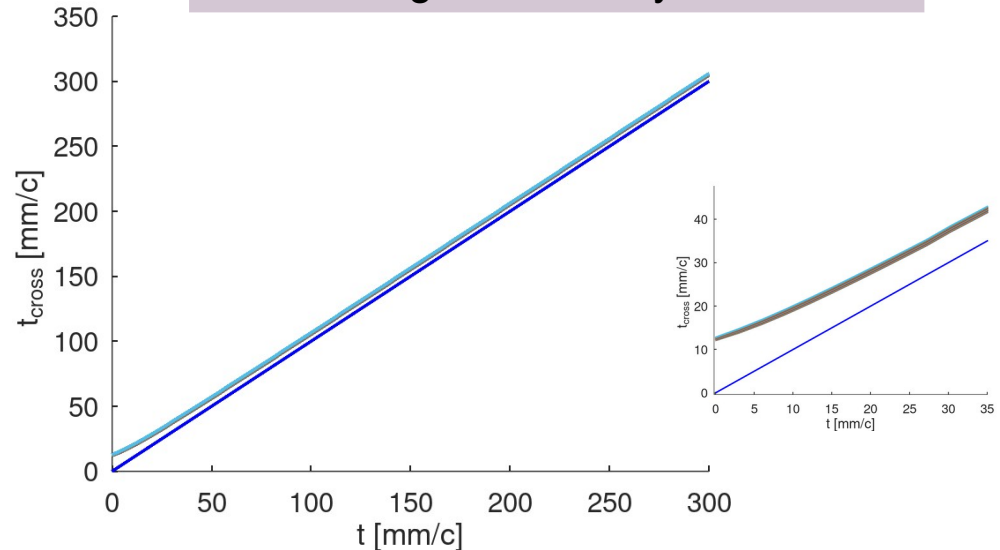
> CLEAR photo-injector information [5].

[5] <https://clear.cern/content/beam-line-description>

$$\exists t^* \in \mathbb{R} \text{ so that } t^* = \frac{a}{c} + \frac{1}{c} \sqrt{a^2 + (s_0 + z(t^*))^2}$$

Catch-up condition

No crossing → Causality not violated



> Catch-up condition verification plot with s_0 ranging from 0 to 3σ .

PART III: BL Simulations in RF-Track

RF-Track

- About **RF-Track** [6]:
 - Beam tracking in field maps/analytic structures including **space-charge** effects, **wakefields**, ...
 - Multiple species (arbitrary q and m)
 - **Parallel C++**, interface with user via **Octave** or **Python**
- **Beam Loading** in RF-Track:
 - Self-consistent module
 - Additional decelerating kick (F_{BL})
 - Attached to Drift spaces, Analytic TW & SW structures, field maps

[6] A. Latina. *RF-Track Reference Manual*. CERN, Geneva, Switzerland, June 2020 DOI: 10.5281/zenodo.3887085

BL in RF-Track

- Based on numerical resolution of $\frac{\partial G_{\text{eff}}}{\partial t} = -\frac{\omega}{Q} \frac{G_{\text{eff}}}{2} - \frac{\omega r_{\text{eff}} \tilde{I}}{2Q} + \frac{G_{\text{target}} \omega}{2Q}$
 - Finite difference method

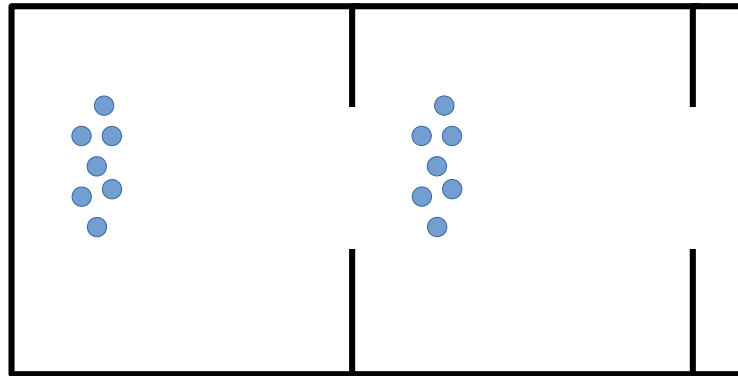
$$t \rightarrow \{t_m\}_{m=0}^{M-1} \quad z \rightarrow \{z_n\}_{n=0}^{N-1} \quad G(z, t) \rightarrow G(z_n, t_m) := G(n, m) \quad \tilde{I}(z, t) \rightarrow \tilde{I}(n, m)$$

$$G(n, m+1) = \left(1 - \frac{\omega \Delta t}{2Q}\right) G(n, m) - C(n, m, \beta) + \frac{G_{\text{target}} \omega}{2Q}$$

$$C(n, m, \beta) = \frac{\Delta t \omega r / Q \mathcal{T}(n, m, \beta) \tilde{I}(n, m, \beta)}{2}$$

BL in RF-Track

- **INPUT:** BEAM, Q , r/Q , $E_z(z, t=0)$, t_{inj}
- **Phase 1:** Preparation: $G(z, t=0)$; Initialize $m = 0$
- **Phase 2:** Compute force while tracking



> Tracking simulation sketch.

At $t_m = m\Delta t$

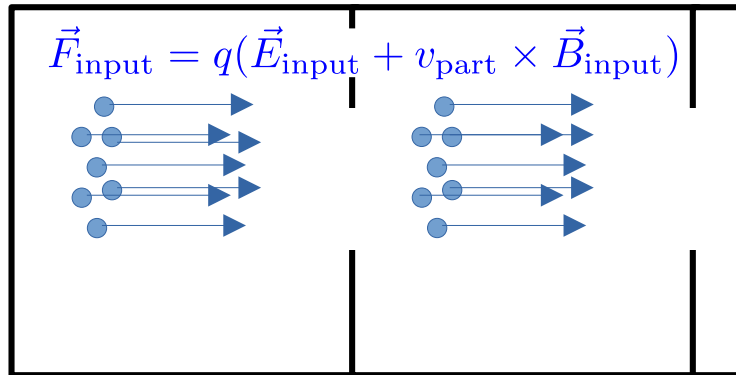
$\tilde{I}(z_n, t_m)$

Get $G(n, m + 1)$

$$\text{Define } \chi(n, m) = 1 - \frac{G(n, m + 1)}{G(n, 0)}$$

BL in RF-Track

- **INPUT:** BEAM, Q , r/Q , $E_z(z, t=0)$, t_{inj}
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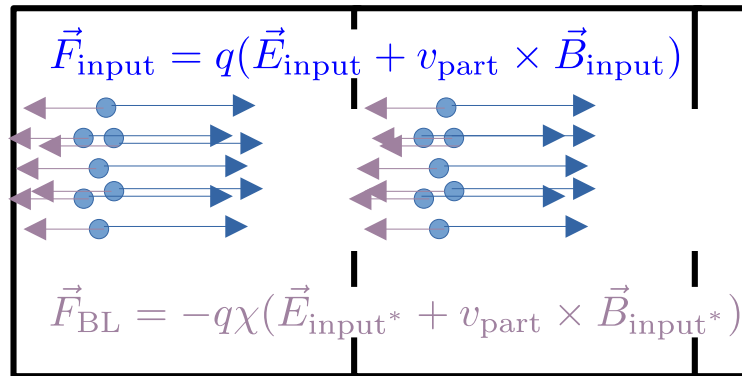
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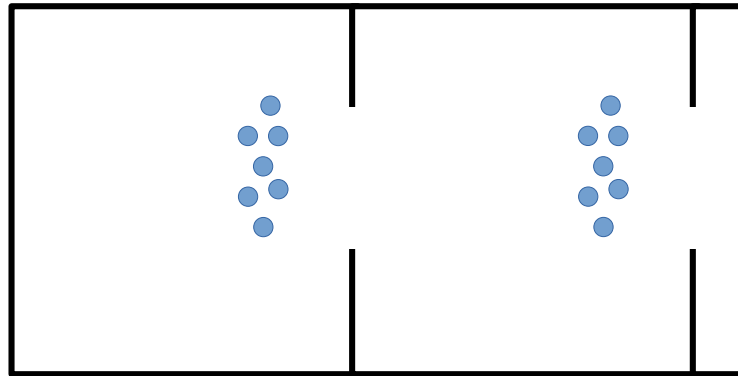
$\tilde{I}(z_n, t_m)$

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BL in RF-Track

- **INPUT:** BEAM, Q , r/Q , $E_z(z, t=0)$, t_{inj}
- **Phase 1:** Preparation: $G(z, t=0)$; Initialize $m = 0$
- **Phase 2:** Compute force while tracking



> Tracking simulation sketch.

At $t_{m+1} = (m + 1)\Delta t$

$\tilde{I}(z_n, t_{m+1})$

Get $G(n, m + 2)$

Define $\chi(n, m + 1) = 1 - \frac{G(n, m + 2)}{G(n, 0)}$

Stop when all particles leave volume

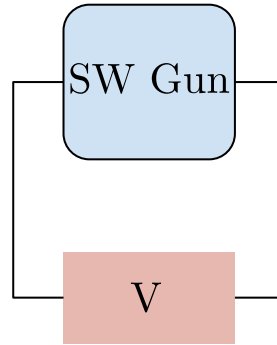
PART IV: Results

BL @ CLEAR photo-injector

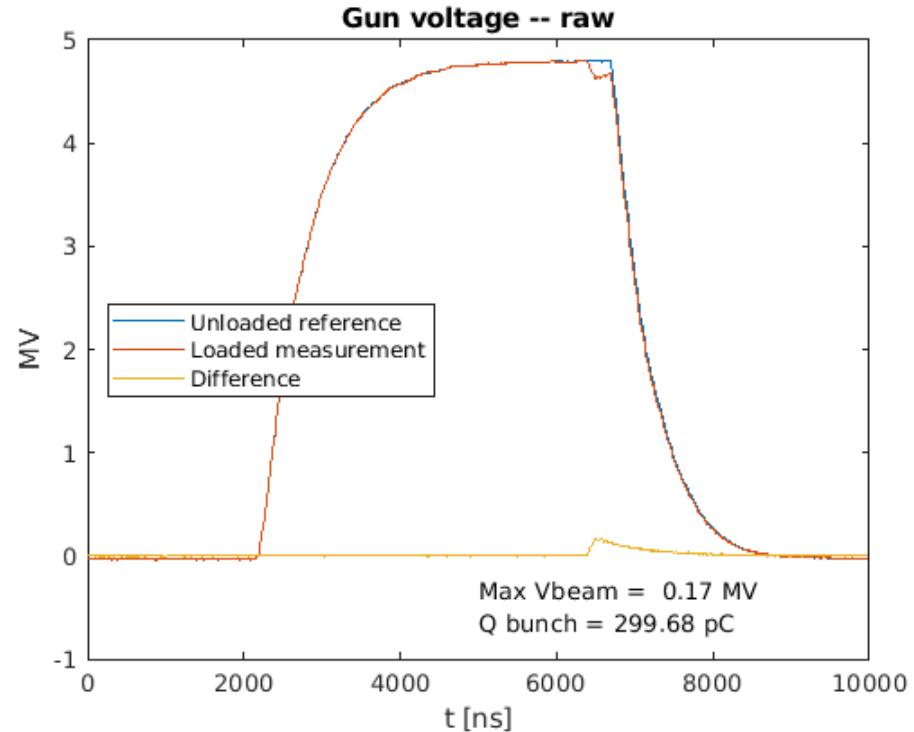
- Voltage measurements for gun-loop

Magnitude	Units	Value
f	GHz	2.997
ϕ_{ad}	rad	π
L_{total}	m	0.175
Q_0		15773
Q_l		5920
r/Q	Ω/m	3080
E_z^{max}	MV/m	60.8

> CLEAR photo-injector information [5].



> Sketch of voltage measurement at CLEAR photo-injector.



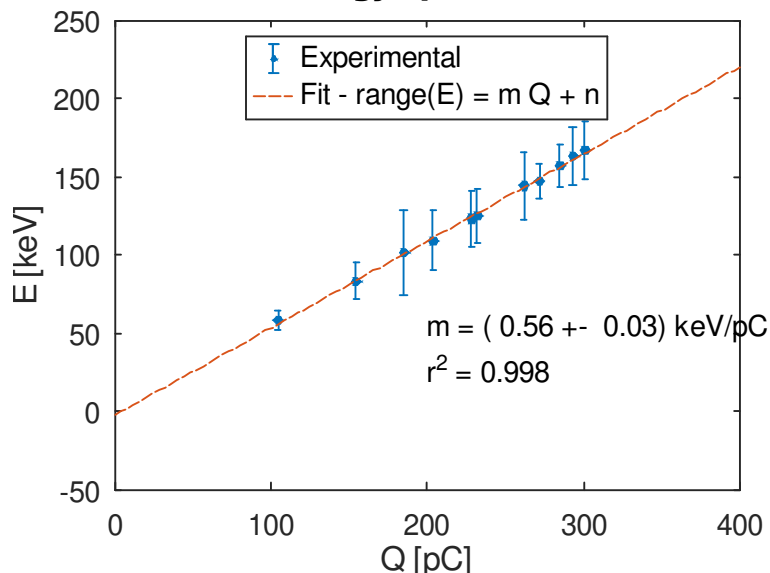
> Voltage along an RF-cycle for CLEAR SW photo-injector. Measurement.

[5] <https://clear.cern/content/beam-line-description>

BL @ CLEAR photo-injector

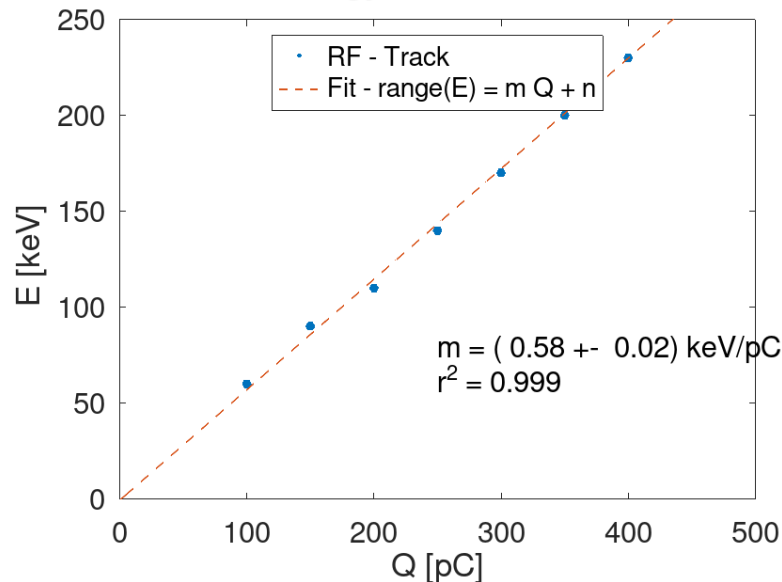
- Train of 150 bunches with variable charge (Q_{bunch}) per bunch; $f_b = f_{\text{RF}}/2$

BL induced Energy Spread for 150 bunches



> Beam Loading Energy Spread induced in a train of 150 bunches as a function of charge

RF-Track Energy loss for the #150 bunch



> Beam Loading Energy Spread induced in a train of 150 bunches as a function of charge (RF-Track, corrected attempt)

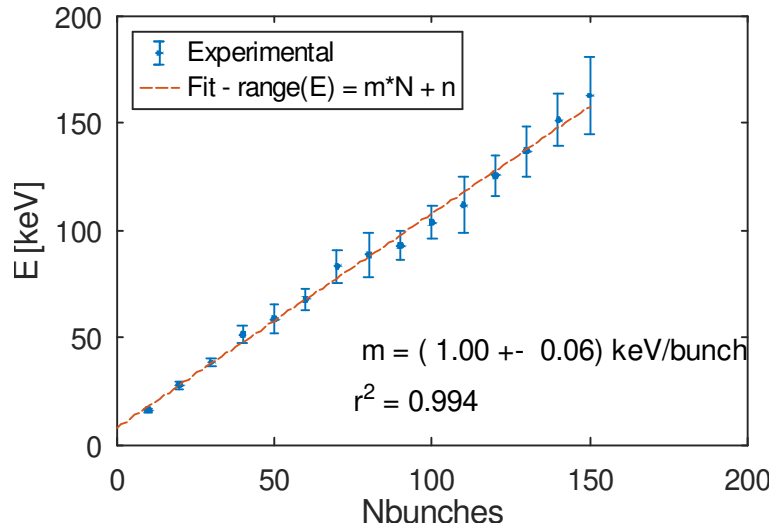
$$\delta_m = \frac{|m_{\text{exp}} - m_{\text{RF-Track}}|}{m_{\text{exp}}}$$

$$\delta_m = 3.5\%$$

BL @ CLEAR photo-injector

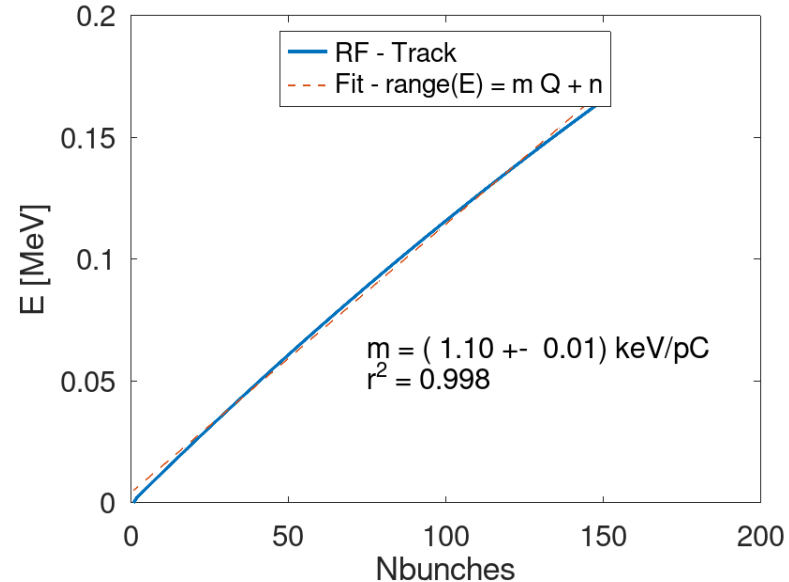
- Train of variable N_{bunches} with fixed charged per bunch

BL induced Energy Spread vs Nbunches
Q bunch = (294.06 ± 25.33) pC



> Beam Loading Energy Spread induced in a train with varying number of bunches and fixed charge per bunch.

BL induced Energy Spread ($Q_{\text{bunch}} = 294 \text{ pC}$)

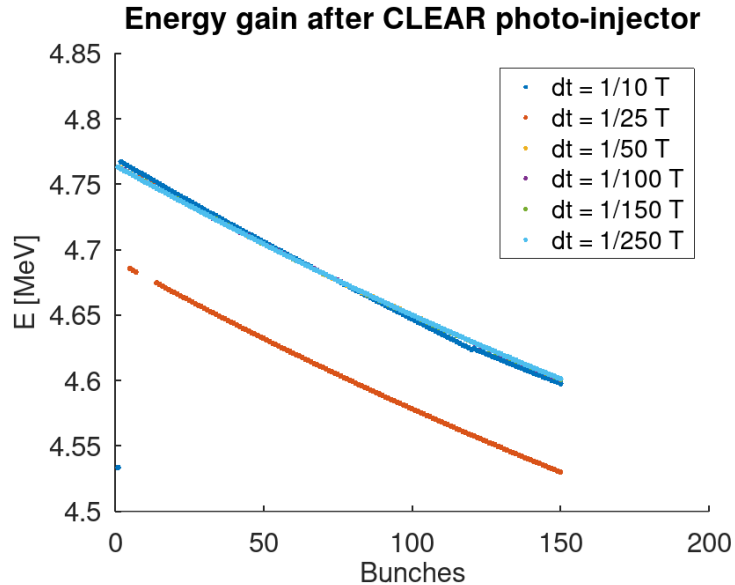


> Beam Loading Energy Spread induced in a train with varying number of bunches and fixed charge per bunch.

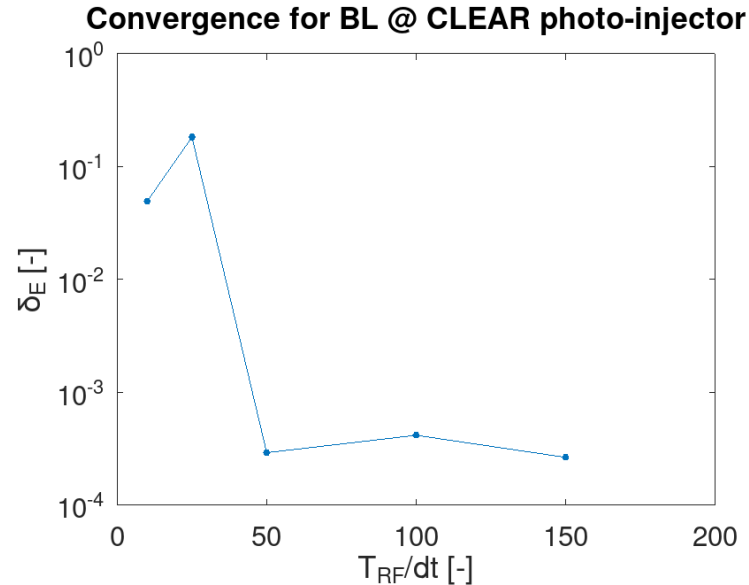
$$\delta_m = 10.0\%$$

BL in RF-Track. Convergence

- Proof of **convergence** under suitable ϕ/t choice



> Energy gain of a train of 150 bunches with a charge per bunch of 300 pC.



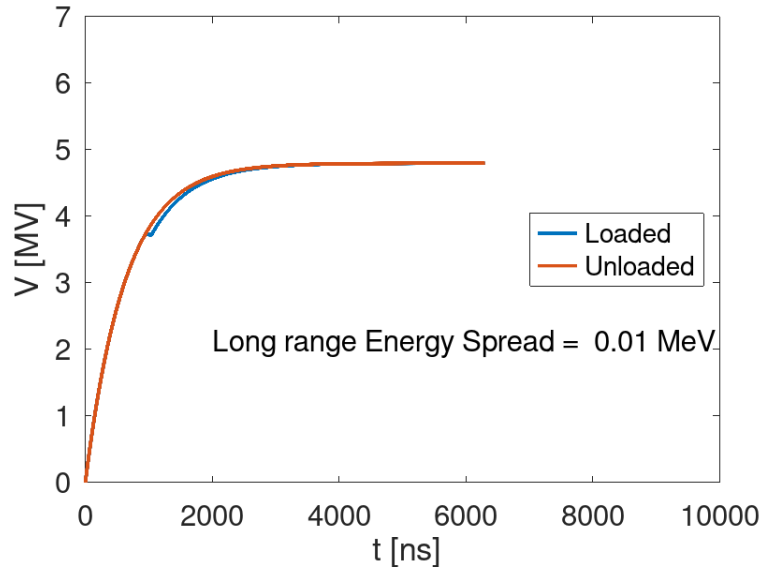
> Relative deviation of the train energy with respect to a highly accurate (E1000) energy distribution as a function of the dt/T fraction.

$$\delta_{E,n} = \sqrt{\sum \frac{(E_n - E_{1000})^2}{E_{1000}^2}}$$

BL compensation

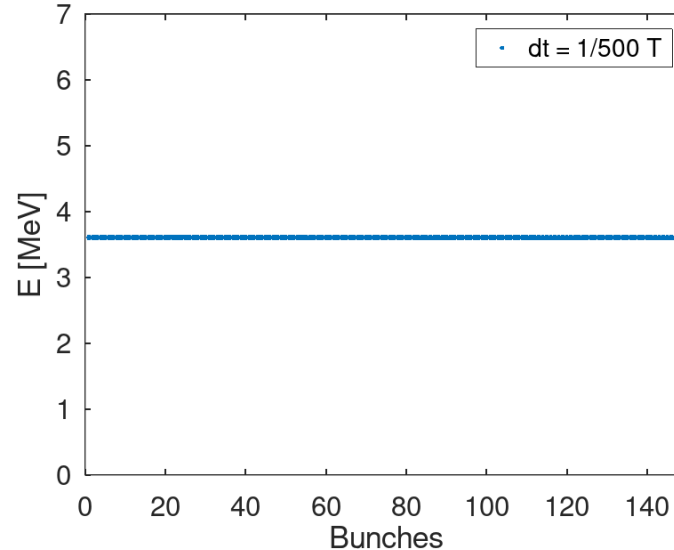
- BL can be compensated with **early injection** of the particles
 - **RF-Track** allows the simulation of this scenario

Voltage cycle for the injector. $t_i/t_f = 1.5$



> Voltage along an RF-cycle for CLEAR SW photo-injector.

Energy gain after CLEAR photo-injector with $t/t_{inj} = 1.5$



> Energy gain of a train of 150 bunches with a charge per bunch of 300 pC.

$$t_{inj} = \frac{2Q_l}{\omega}$$

Conclusions

- Beam Loading effect causes gradient reduction in **high-gradient & high-intensity** accelerators.
- The implementation of this model in **RF-Track** provides a user-friendly, flexible and powerful tool which shows:
 - Numerical **convergence & stability**

... and allows:

- Simulation of **realistic** scenarios
 - Good **agreement with measurements** at the CLEAR photo-injector
- Beam Loading **compensation**

OUTLOOK: Further development of the code: Interaction of BL with other collective-effects & Start-to-end simulations.

Acknowledgements

- **Supervision**, guidance and trust:
 - Andrea Latina (CERN, BE-ABP-LAF)
also, RF-Track creator and main developer.
 - Nuria Fúster Martínez, Benito Gimeno, Daniel Esperante (UV – Spain, IFIC - CSIC)
- Useful **material & discussions**:
 - Alexej Grudiev (CERN, SY-RF-MKS)
 - Avni Aksoy (CERN, BE-ABP-LAF)
 - The CLEAR OP team

Thanks for your attention

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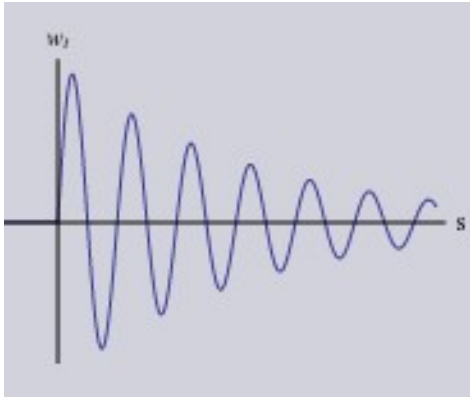
BACK UP SLIDES

Back up slide – Causality (General)

- **Shunt impedance** is defined as:

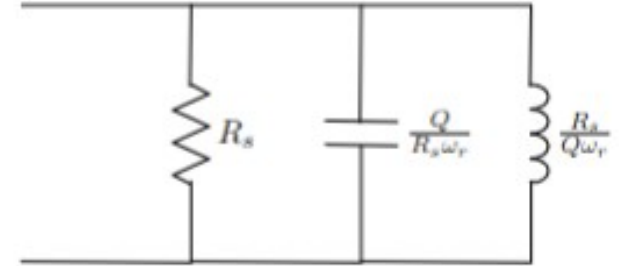
$$R_s = \frac{V_{\text{acc}}^2}{P_{\text{diss}}}$$

- Such definition comes from the LCR circuit analogy
 - Which comes from the fact of considering a **causal longitudinal wakefield**



Causal longitudinal wake

$$\implies Z_l \simeq \frac{R}{1 + jQ \left(\frac{\omega_r}{\omega} - \frac{\omega}{\omega_r} \right)} \implies$$



Accelerator behaves as parallel LCR circuit in resistive mode

GUN: (E_z , φ) Calibration

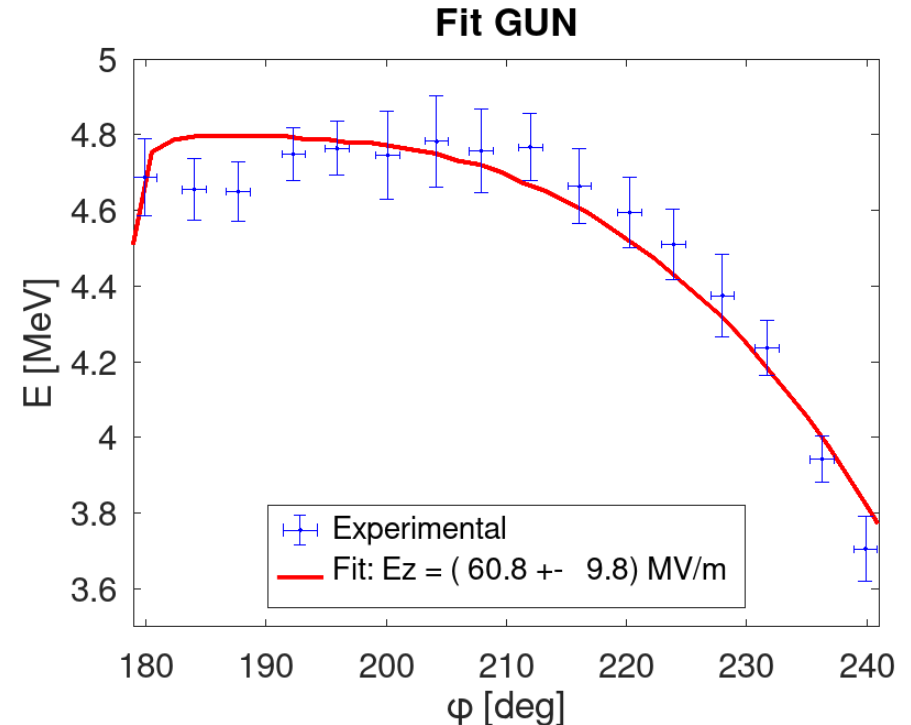
With Avni Aksoy

- For a 50 pC bunch:
 - Collect (E_k , φ_k) measurements
 - Fit then to target function $F(E_z, \varphi)$
 - F: RF-Track calculation of E after gun.

Magnitude	Units	Value
E_z^{\max}	MV/m	60.8 ± 9.8
r^2		0.94

> Results of the minimum square fitting with a test function computed with RF-Track.

Still cathode influence to be studied!



> Energy gain after the gun as a function of the phase. In red, RF-Track prediction. In blue, experimental results

GUN: Beam Loading Measurements

- 2) Divergent slope

- Looking again at the BL equation ...

$$-\frac{\partial G}{\partial t} = \frac{\omega}{Q} \frac{G}{2} + \frac{\omega r_{\text{eff}} I}{2Q} + \frac{G_{\text{init}} \omega}{2Q}$$

... the slope of the plot E vs Q depends on **r/Q** and **Q**

- From design report: $r/Q =$; $Q_0 = 14530$;
- However, we learn that the Q governing the dynamics is

$$Q = (598 \pm 8) \cdot 10$$

- This is the loaded quality factor!

$$Q_l = \frac{Q_0}{1 + \beta} = (598 \pm 8) \cdot 10 \implies \beta = 1.5$$

