

# Slice Energy Spread Measurements in the EuXFEL Injector

**Stuart Walker**, Sergey Tomin, Igor Zagorodnov, Nina Golubeva, Matthias Scholz, Winfried Decking, Bolko Beutner, Erion Gjonaj

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# Overview

## Subheading, optional

**1 Motivation**

**2 Context**

**3 Overview of the EuXFEL**

**4 Method**

**5 Measurements with Results**

**6 Conclusion**

# Motivation

- Uncorrelated (slice) energy spread is an important property in FELs, where high brightness is required.
- However, usually energy spread from injector is too low and must be increased using the laser heater (LH) to increase SASE performance due to the microbunching instability.
  - Ideally increased to 8 keV based on simulations from Martin Dohlus, but actual optimal value never measured.
- Regardless, understanding uncorrelated energy spread is necessary to improve machine reproducibility and understanding.
- Dynamics not understood, unable to recreate measured slice energy spread values with ASTRA simulation.

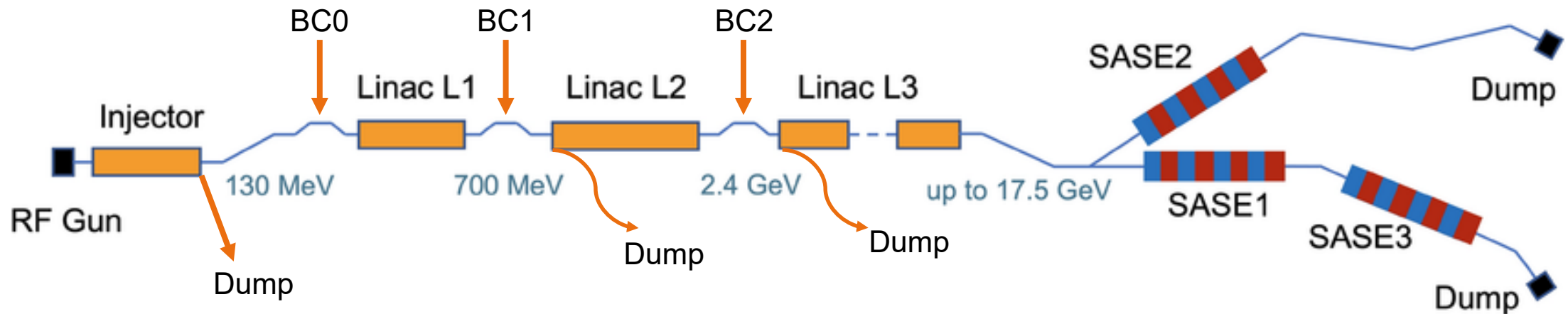
# Context

- Fermi, 2020: Intrabeam scattering measured, but result highly dependent on initial energy spread, which was treated as a free parameter.
- SwissFEL, 2020: energy spread measurement in the injector using an energy scan—15 keV @ 200pC & >100MeV.
- EuXFEL, 2021: Similar approaching involving a dispersion scan—6 keV @ 250pC & 130MeV.
- PITZ, 2022: 2 keV @ 250pC & 20MeV.
- SwissFEL, 2022: Contribution of the microbunching instability and intra-beam scattering evidenced by varying the optics -> reduced by adjusting optics and disabling LH chicane.
- This talk: recent measurements and new results since November 2022 in the EuXFEL injector.

# The EuXFEL

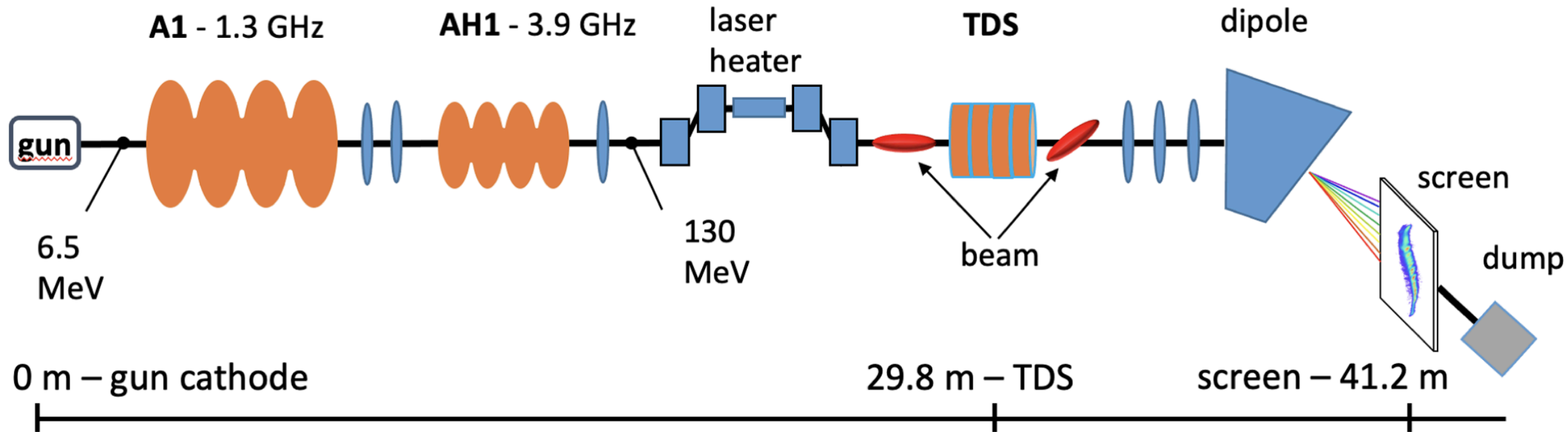
## Overview

- 3.1 km machine.
- Two hard x-ray undulator lines, SA1 and SA2.
- One soft x-ray line, SA3.
- Four chicanes, laser heater chicane in injector and three for compression.
- Diagnostic stations in injector and after BC2 (at max compression).
- One transverse deflecting structure (TDS) in the injector, and one after BC2.
- Beam can be matched in diagnostic sections and just after BC1.
- Hard x-ray self-seeding in SA2, wish recent pushes to more special modes.
- Crucial to know longitudinal beam dynamics to deliver the best performance.



# The EuXFEL

## Injector



# The Usual Approach

## Energy Spread Measurement

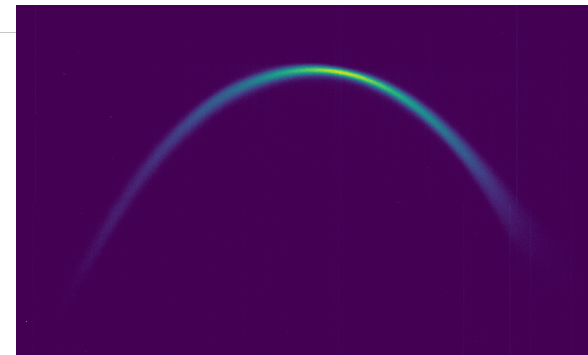
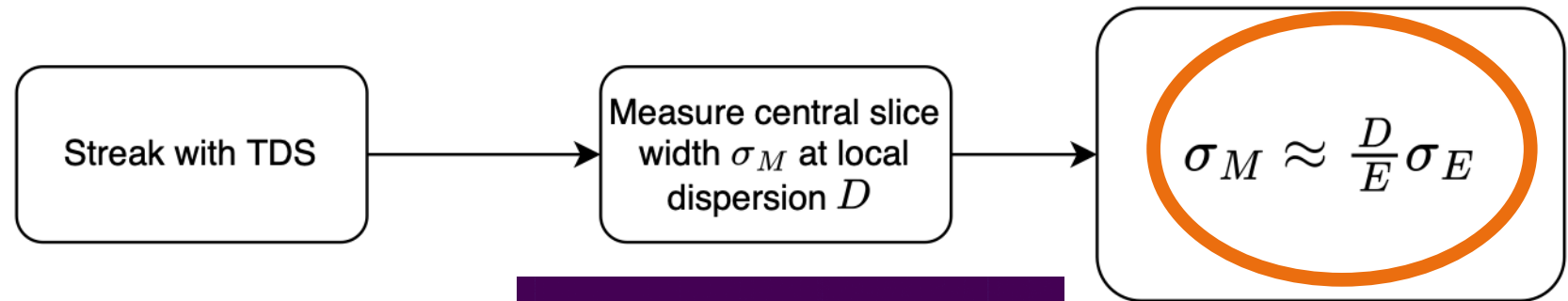
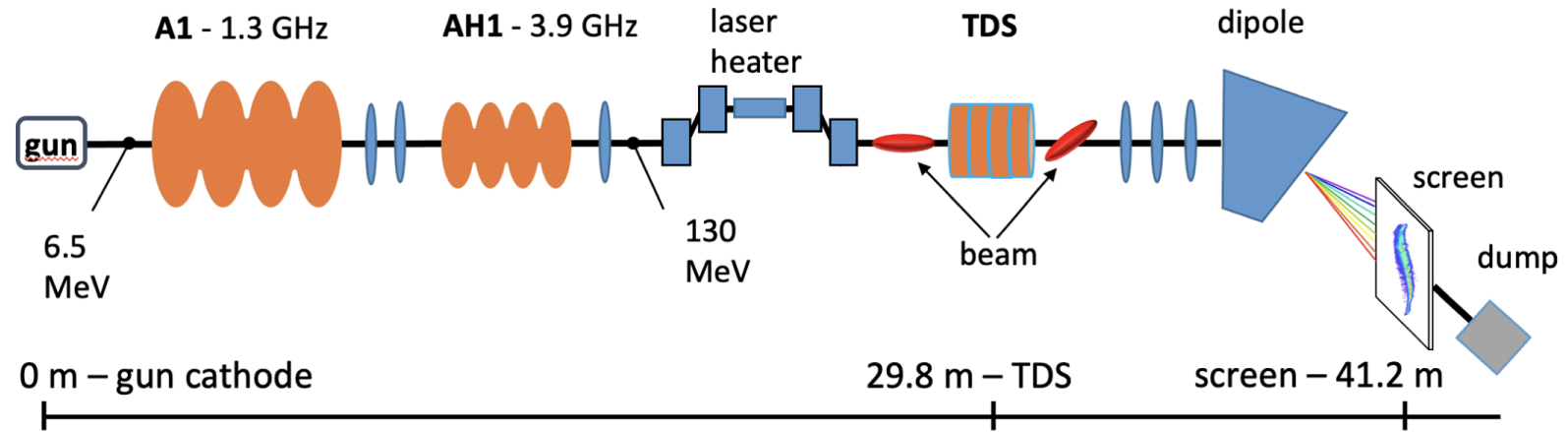
Only provides an upper limit on uncorrelated spread  $\sigma_E$ .

Neglected contributions to size:

- Intrinsic betatronic beam size
- TDS-induced energy spread
- Imaging resolution

In the injector these contributions can be larger than the slice energy spread's contribution!

Need to separate these effects...



# Separating the contributions to the slice size

## Energy Spread Measurement

Assuming no correlation between TDS-induced energy spread and "true" energy spread, the final slice energy spread seen at the screen:

$$\sigma_{E_{\text{final}}}^2 = \sigma_E^2 + (ekV)^2 \sigma_I^2$$

Beamsize in the TDS

$$\sigma_M^2 = \sigma_R^2 + \sigma_B^2 + \frac{D^2}{E^2} \sigma_E^2 + \frac{D^2}{E^2} (ekV)^2 \sigma_I^2$$

Imaging resolution

Betatron

energy spread contribution

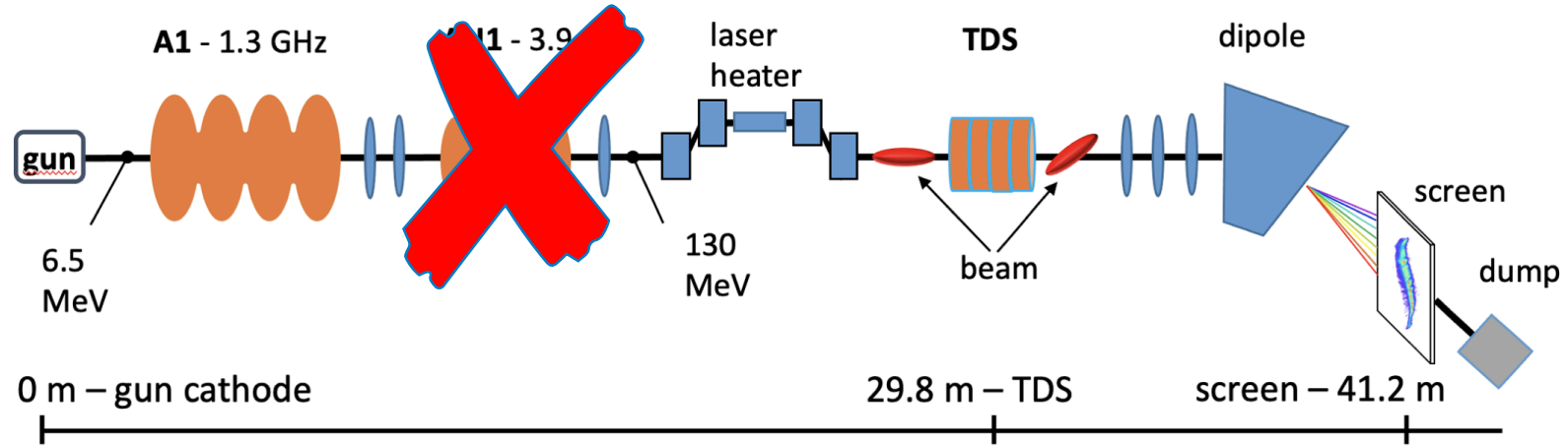
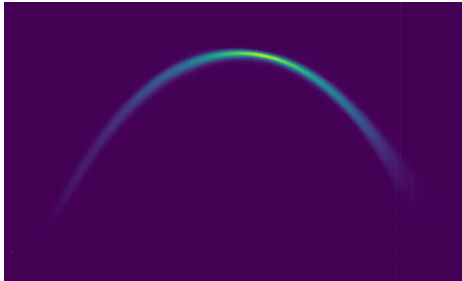
TDS contribution

$$\sigma_B^2 = \frac{\beta_x \varepsilon_n}{\gamma_0} \quad \sigma_I^2 = \frac{\varepsilon_n}{\gamma_0} (\beta_y^0 + 0.25 L^2 \gamma_y^0 - L \alpha_y^0)$$

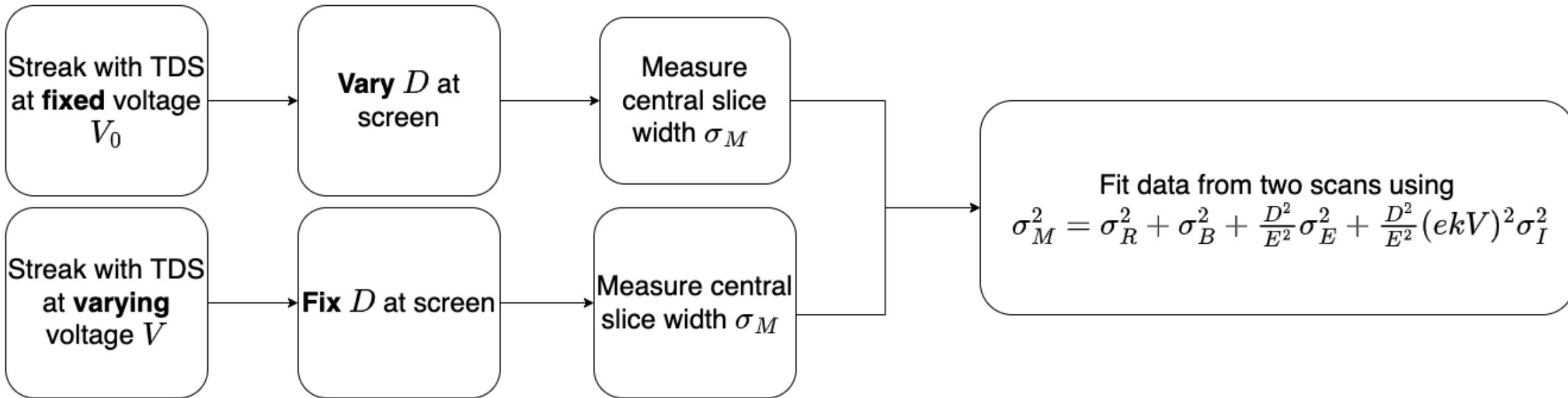


# High Resolution Technique

## Energy Spread Measurement



Dispersion Scan  
TDS Scan



# Derived Values

## Energy Spread Measurement

For the two scans:

$$\sigma_M^2 = \sigma_R^2 + \sigma_B^2 + \frac{D^2}{E^2} \sigma_E^2 + \frac{D^2}{E^2} (ekV)^2 \sigma_I^2$$

$$\sigma_M^2 = A_V + B_V V^2 \quad \sigma_M^2 = A_D + B_D D^2$$

$$A_V = \sigma_R^2 + \sigma_B^2 + \frac{D^2}{E^2} \sigma_E^2$$

$$A_D = \sigma_R^2 + \sigma_B^2$$

$$B_V = \frac{D^2}{E^2} (ek)^2 \sigma_I^2$$

$$B_D = \frac{1}{E^2} \sigma_E^2 + \frac{1}{E^2} (ekV)^2 \sigma_I^2$$

$$\sigma_E = \frac{E_0}{D_0} \sqrt{A_V - A_D}$$

$$\sigma_I = \frac{E_0}{D_0 ek} \sqrt{B_V}$$

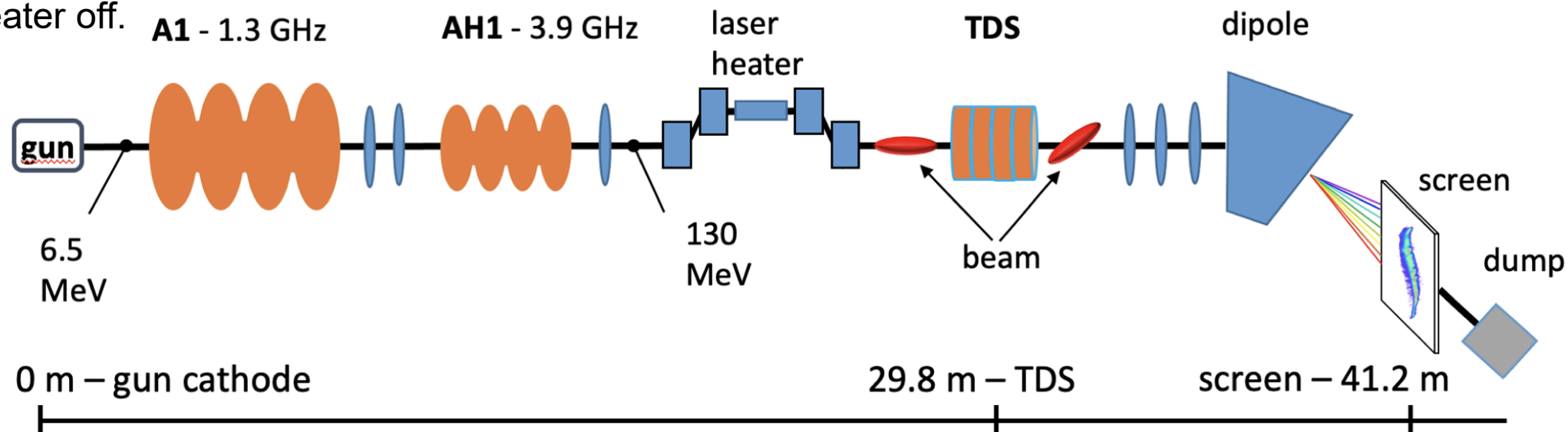
$$\sigma_B = \sqrt{B_\beta \beta_x^0}$$

$$B_\beta = \sigma_I^2 (\beta_y^0 + 0.25L^2 \gamma_y^0 - L\alpha_y^0)^{-1}$$

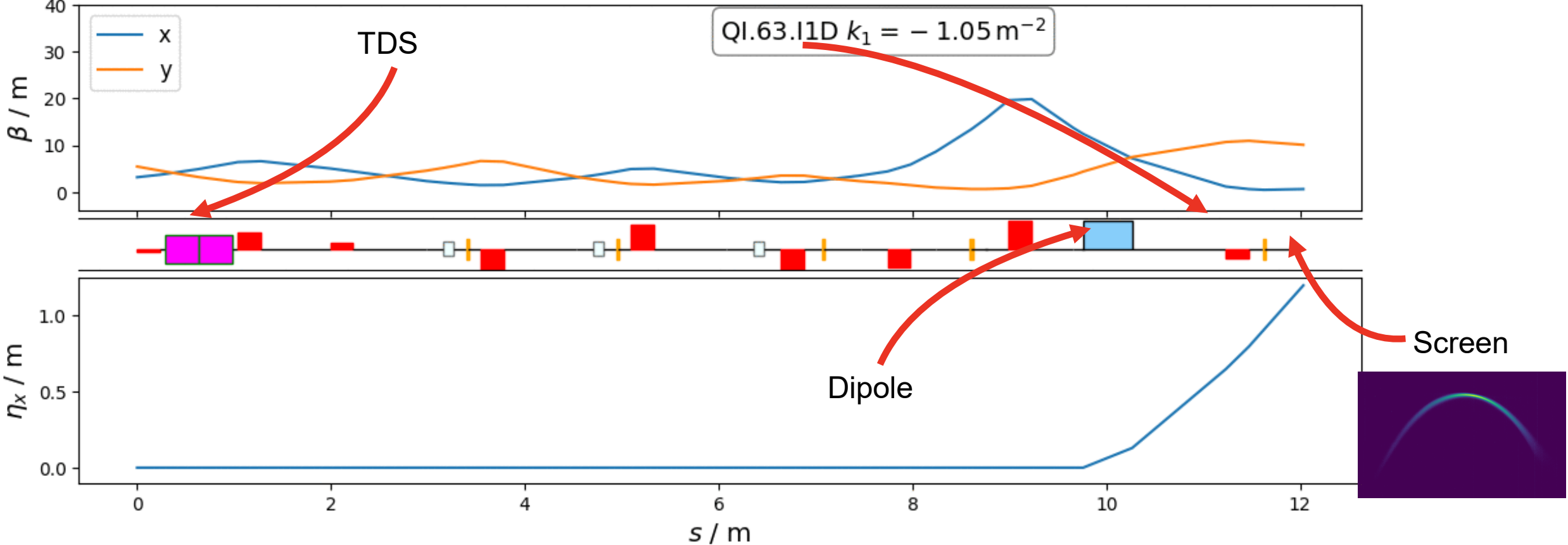
$$\sigma_R = \sqrt{A_D - \sigma_B^2}$$

# Injector Setup

- Calibrate the gun phase.
- Turn off AH1 for minimum chirp contribution to energy spread.
- Go on crest in A1, adjust the voltage so we are at 130 MeV at the screen.
- Turn the laser heater off.
- Apply **special optics** to maximise the ratio of the dispersion to the betatron contributions to the spot size.
- Match the central slice.
- Measure the dispersion.



# Special optics for measurement from TDS to the screen



# Procedure

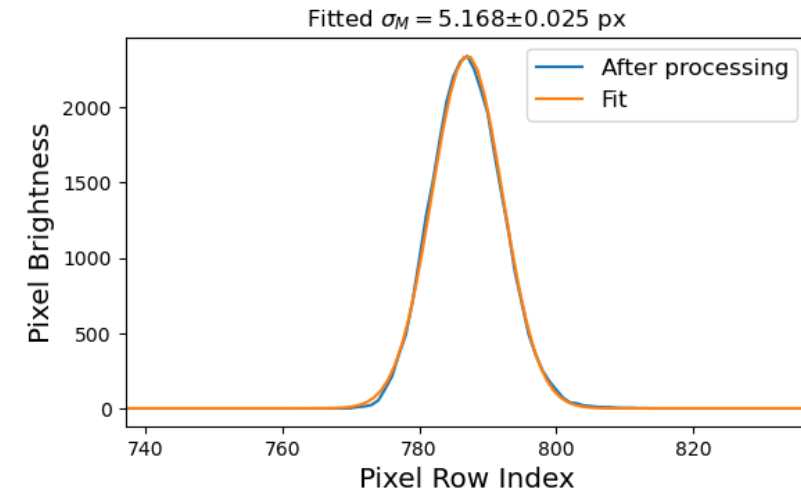
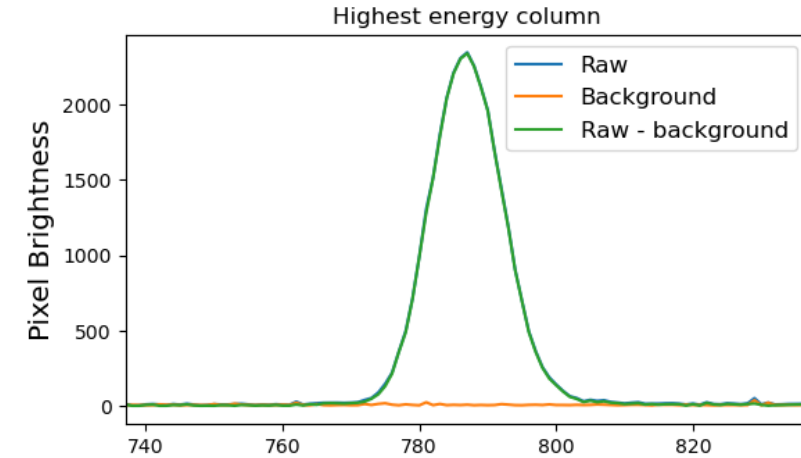
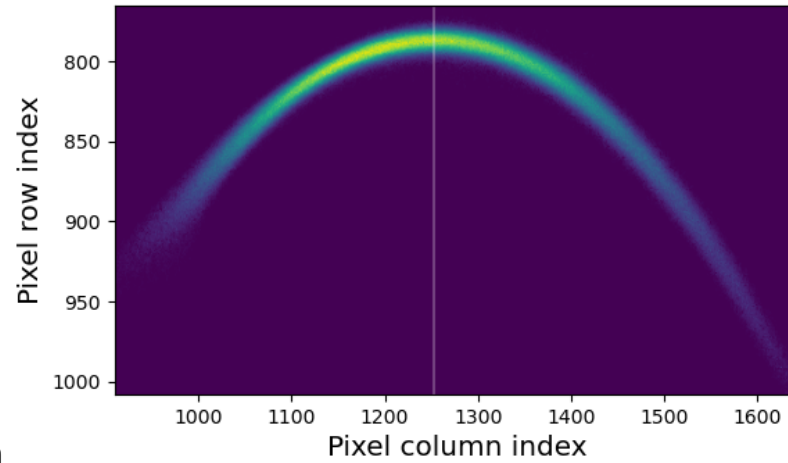
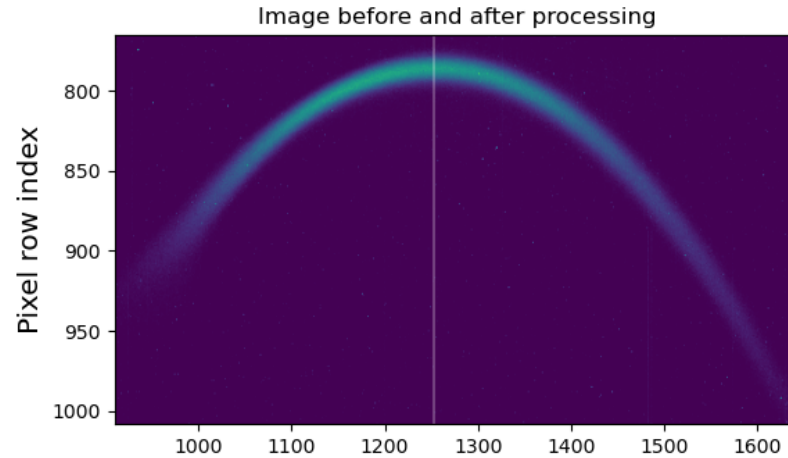
TDS No. = 13,  $\eta_x = 1.181$  m, image 1, before/after image processing

Take 5 background images at the start of the measurement, and then 30 images:

1. At each TDS  $V$  in the voltage scan.
2. At each  $D_x$  at the screen in dispersion scan.
3. At each  $\beta_x$  at the screen in the beta scan.

Then:

1. Subtract background.
2. Mask to remove isolated blobs.
3. Pick the largest connected non-zero pixel blob to be "the beam".
4. Fit a Gaussian to 10 slices centred on highest energy slice & average.



# 28th November 2022 Measurements

- Aim:

- Repeat measurement from 2021 publication ( $\sigma_E = 5.9$  keV).
- Match beamline configuration from that time as much as possible.
- Particular attention to:
  - Gun gradient: 54.7 MV/m.
  - Solenoid: 326 A.
- **Cathode changed in 2022!**

- Outcome:

- Three energy spread measurements over 8 hours:
  1. Gun gradient 54.7 MV/m, solenoid 326 A, gun phase  $-43^\circ$  (most similar to 2021 result).
  2. Gun gradient to 56.5 MV/m, solenoid to 336 A.
  3. Solenoid to 335 A, gun phase to  $-41^\circ$ .

# Results

Scenario	$\sigma_E / \text{keV}$	Solenoid / A	Gun Gradient / $\text{MV m}^{-1}$	Gun Phase / $^\circ$
Feb. 2021 Published Result	5.842	338	56.7	-42.9
Nov. 2022 Measurement 1	4.313	326.6	54.7	-43.1
Nov. 2022 Measurement 2	3.635	336	56.5	-43.6
Nov. 2022 Measurement 3	3.385	335	56.5	-41.6

Variable	Feb. 2021	Nov. 2022	Unit
$A_V$	$4.530 \times 10^{-9}$	$3.043 \times 10^{-9}$	$\text{m}^2$
$B_V$	$1.515 \times 10^{-21}$	$1.326 \times 10^{-21}$	$\text{m}^2 \text{V}^{-2}$
$A_D$	$1.702 \times 10^{-9}$	$1.540 \times 10^{-9}$	$\text{m}^2$
$B_D$	$2.433 \times 10^{-9}$	$1.416 \times 10^{-9}$	-
$\sigma_E$	5.8	4.3	keV
$\sigma_I$	68.0	64	$\mu\text{m}$
$\sigma_B$	29	28	$\mu\text{m}$
$\sigma_R$	28	27	$\mu\text{m}$
$\varepsilon_n$	0.38	0.34	$\text{mm} \cdot \text{mrad}$
$V_0$	0.61	-0.63	MV
$D_0$	1.183	1.169	m

$$\sigma_M^2 = \sigma_R^2 + \sigma_B^2 + \frac{D^2}{E^2} \sigma_E^2 + \frac{D^2}{E^2} (ekV)^2 \sigma_I^2$$

Nov. 2022 Measurements:

1. Gun gradient 54.7 MV/m, solenoid 326 A, gun phase  $-43^\circ$  (most similar to 2021 result)
2. Gun gradient to 56.5 MV/m, solenoid to 336 A,
3. Solenoid to 335 A, gun phase to  $-41^\circ$ .

Imperfect recreation of 2021 conditions, was as close as we could get.

# Comments

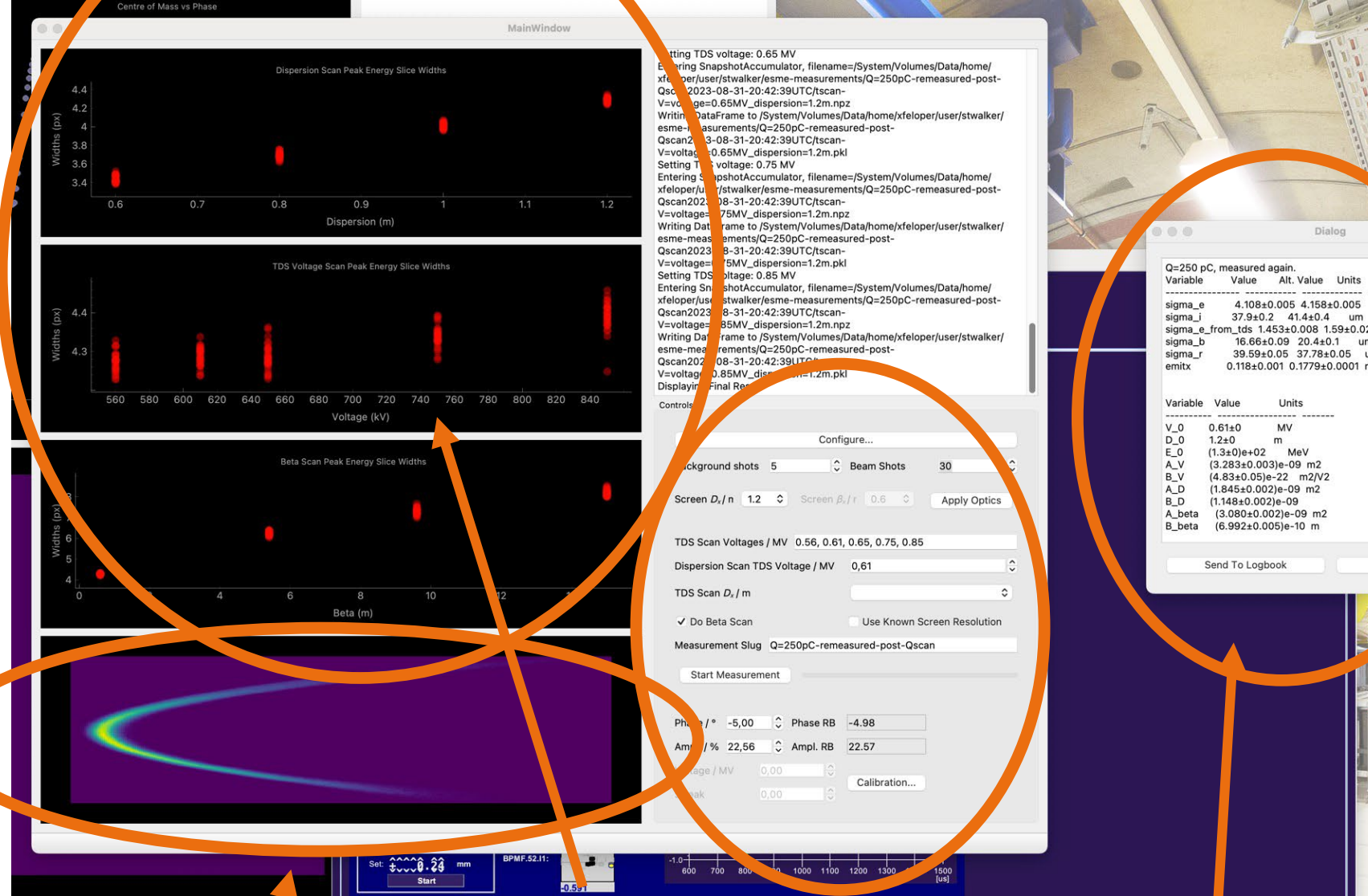
- The measured energy spread in the injector decreased since Feb. 2021 from 5.9 keV to 4.2 keV.
  - Through 2023 has since remained stabled at around 4 keV with repeated measurements.
- Still not the value predicted by theory (Erion's simulations predict  $\approx 3$  keV with IBS).
- In the period between these two measurements:
  - The cathode has changed.
  - The bunch length was increased from around 4 ps to around 5 ps.
- Suggests a long term drift in the energy spread.
- These three datapoints took eight hours to acquire.



# New Tooling

## Measurement

- Taking ~400 images across ~15 machine setpoints without special tooling would take about an hour. Then, offline analysis takes another 10 minutes.
- To this end we have developed a GUI to do the parameter scanning.
- Live image processing panel.
- Live slice width scan plots (e.g. TDS voltage against central slice width).
- Immediate results at the end of the scan.
- **Bottom line: results in a few minutes.**



Post-processed image

Slice widths

Measurement control

Measurement result

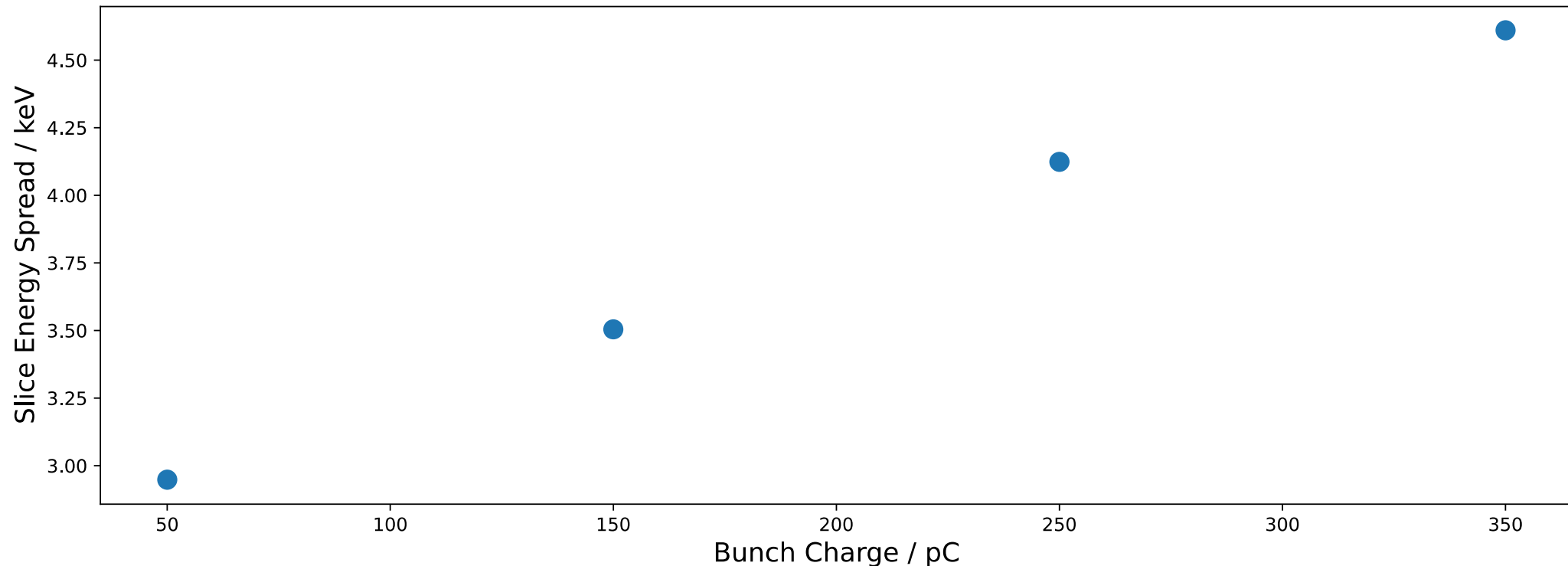
# New Scans

- We did two scans over a single four hour shift on 31<sup>st</sup> August 2023.
  - Bunch charge scan: 50 pC, 150 pC, 250 pC and 350 pC.
  - Laser heater chicane  $R_{56}$  scan: -1 mm, -2.5 mm, -4.336 mm (nominal), -5.5 mm and -7 mm.
- We rematched the projected emittance every single time,
  - then **measured** the central slice emittances (showing good matching already).
- We additionally scanned the beta function at the screen (not shown here).
- The fast measurement software described previously allows us to do these sorts of scans now.

# Bunch Charge Scan

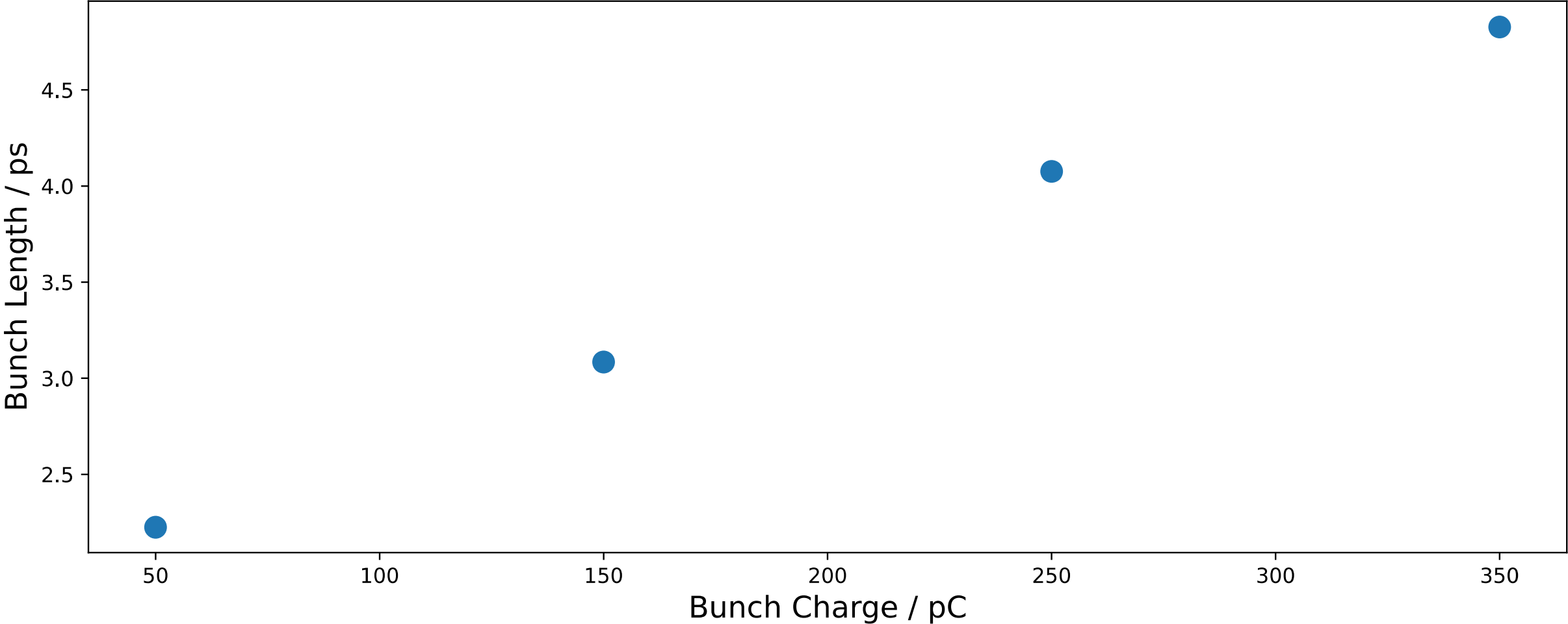
## Slice Energy Spreads

- We scanned from 50 pC to 350 pC in steps of 100 pC.
- Unable to resolve the TDS contribution to the energy spread at 50 pC, so assumed to be zero.
- No obvious reason to consider energy spread from TDS to be zero, related to the camera sensitivity?



# Bunch Charge Scan

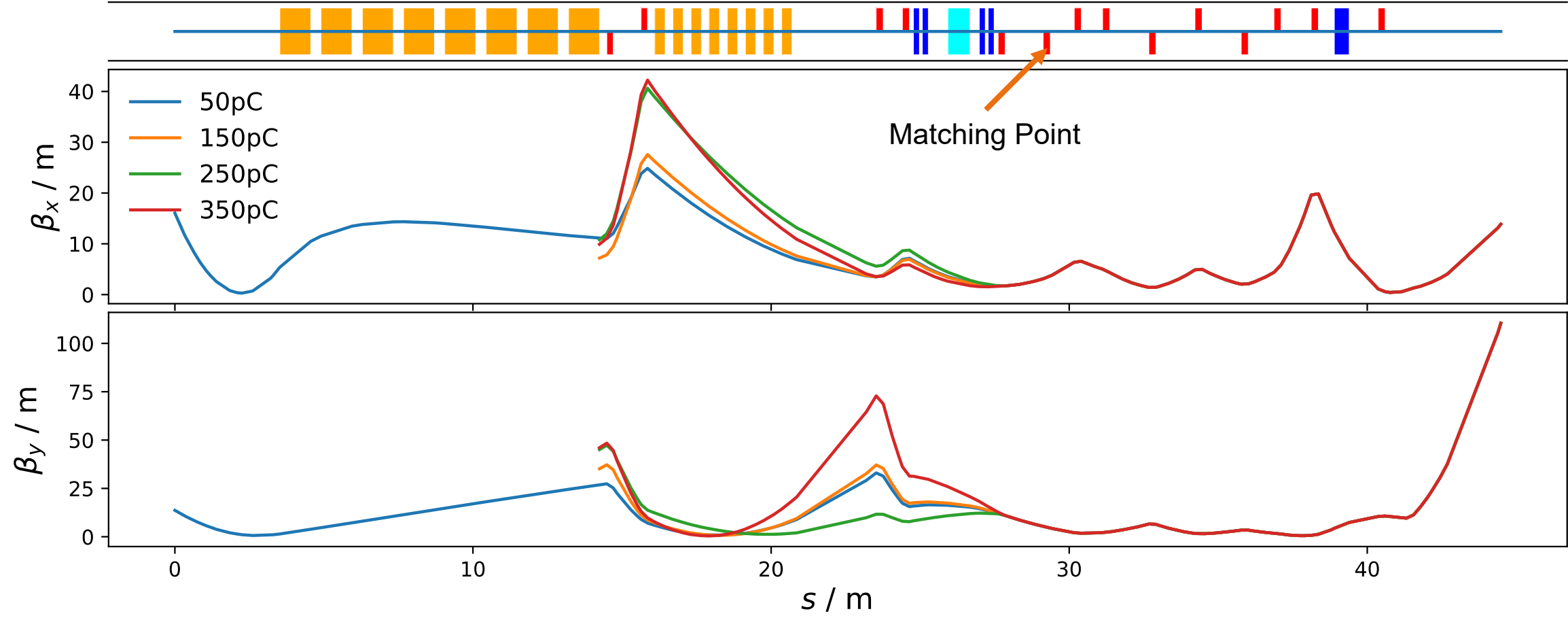
## Bunch Lengths



# Bunch Charge Scan

## Transverse Optics

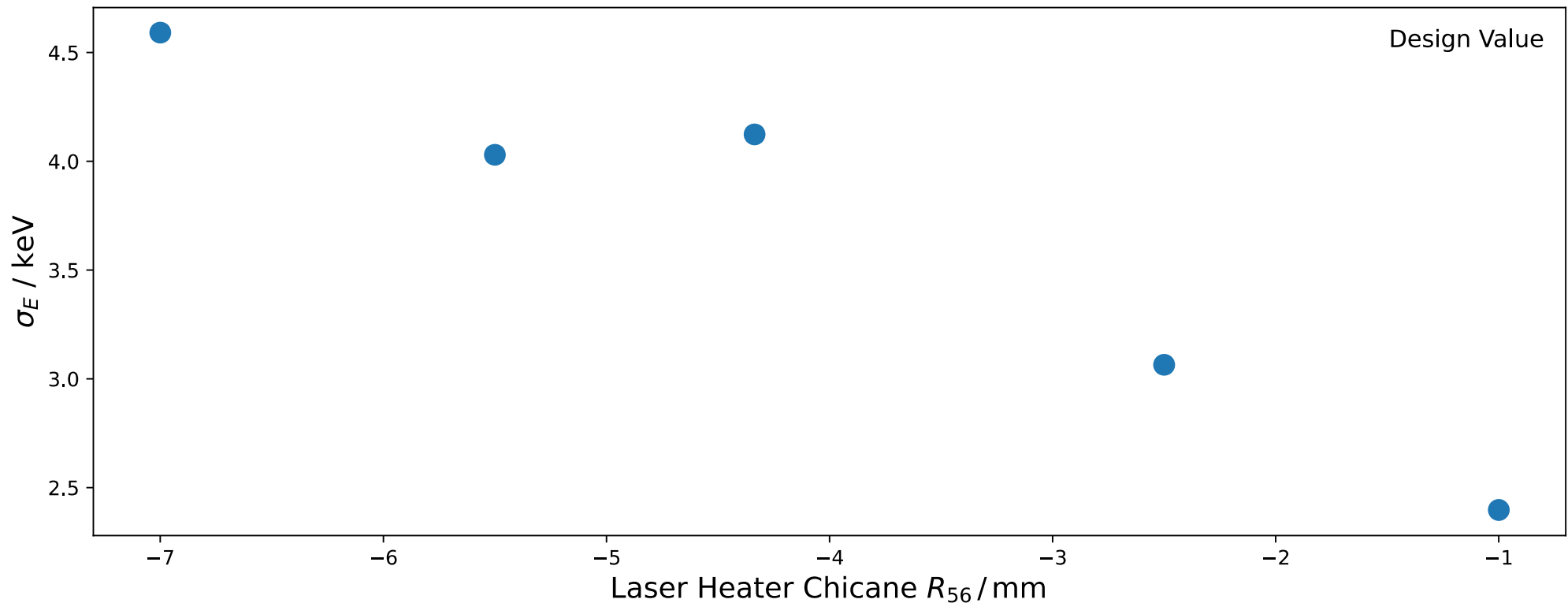
- $D_x = 1.2$  m at the screen.
- Only 50 pC backtracked to cathode due to reduced space charge effects at this charge.



# Laser Heater Chicane $R_{56}$ Scan

## Slice Energy Spreads

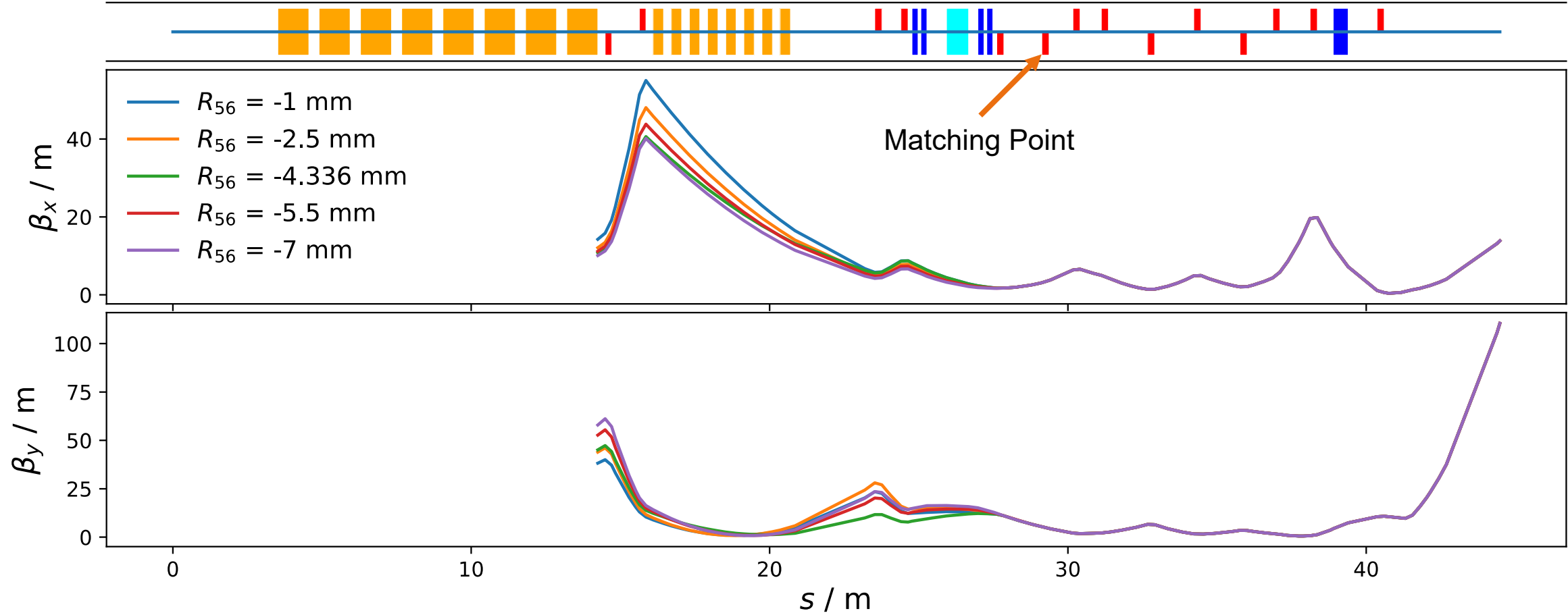
- We scanned the  $R_{56}$  of the laser heater chicane from -7 mm to -1 mm (the minimum) @  $Q = 250$  pC.
- We rematched the beam each time, and then measured the central slice (good matching).
- Design setting (-4.336 mm) not scanned in order with the other setpoints.
- Strong dependence on LH  $R_{56}$  is apparent, from 4.5 keV at maximum to 2.4 keV at minimum.



# Laser Heater Chicane $R_{56}$ Scan

## Transverse Optics

- $D_x = 1.2\text{m}$  at the screen.
- 250 pC, so backtracked only to first cavity.



# Conclusion

- The slice energy spread shows wide tunability based on the charge and laser heater chicane  $R_{56}$ :
  - (2.4–4.6) keV from -1 mm to -7 mm  $R_{56}$  @ 250 pC.
  - (3.0–4.5) keV from 50 pC to 350 pC @ nominal  $R_{56}$ .
- This, in combination with the longer term decrease from 6 keV to 4 keV in the last two years for unknown reasons, makes it clear that there is still a lot to be understood about the longitudinal phase space dynamics in the injector.
- Important to understand to boost reproducibility.
- The charge scan points at possible IBS contributions to the slice energy spread.
- The  $R_{56}$  dependence needs more understanding and explanation. MBI?
- We now have a fast and easy to use data taking and online analysis tool which gives immediate results. Room for further studies of other parameters, e.g. bunch length.



# Outlook

- Upcoming experiment in October for additional chicane studies (opening the undulator and repeating).
- We want to go to 10 pC to neglect collective effects and try to directly observe the thermal slice energy spread from the cathode.
- Aiming to take part in the upcoming new laser commissioning—bunch length scans (not possible with current laser).
- We would like to provide an absolute calibration of the laser heater.
- IBS simulation work ongoing with Erion Gjonaj to ground these results more in theory.
- Further work on the GUI and tooling to improve usability and speed.
- New measurements and studies are possible with new software—open to ideas, suggestions and collaboration!
- We want to ultimately transport to the in BC2 and measure there at 130 MeV—tricky for a number of reasons.

**Thank you**  
**Vielen Dank**

## Contact

Deutsches Elektronen-  
Synchrotron DESY

[www.desy.de](http://www.desy.de)

Stuart Walker

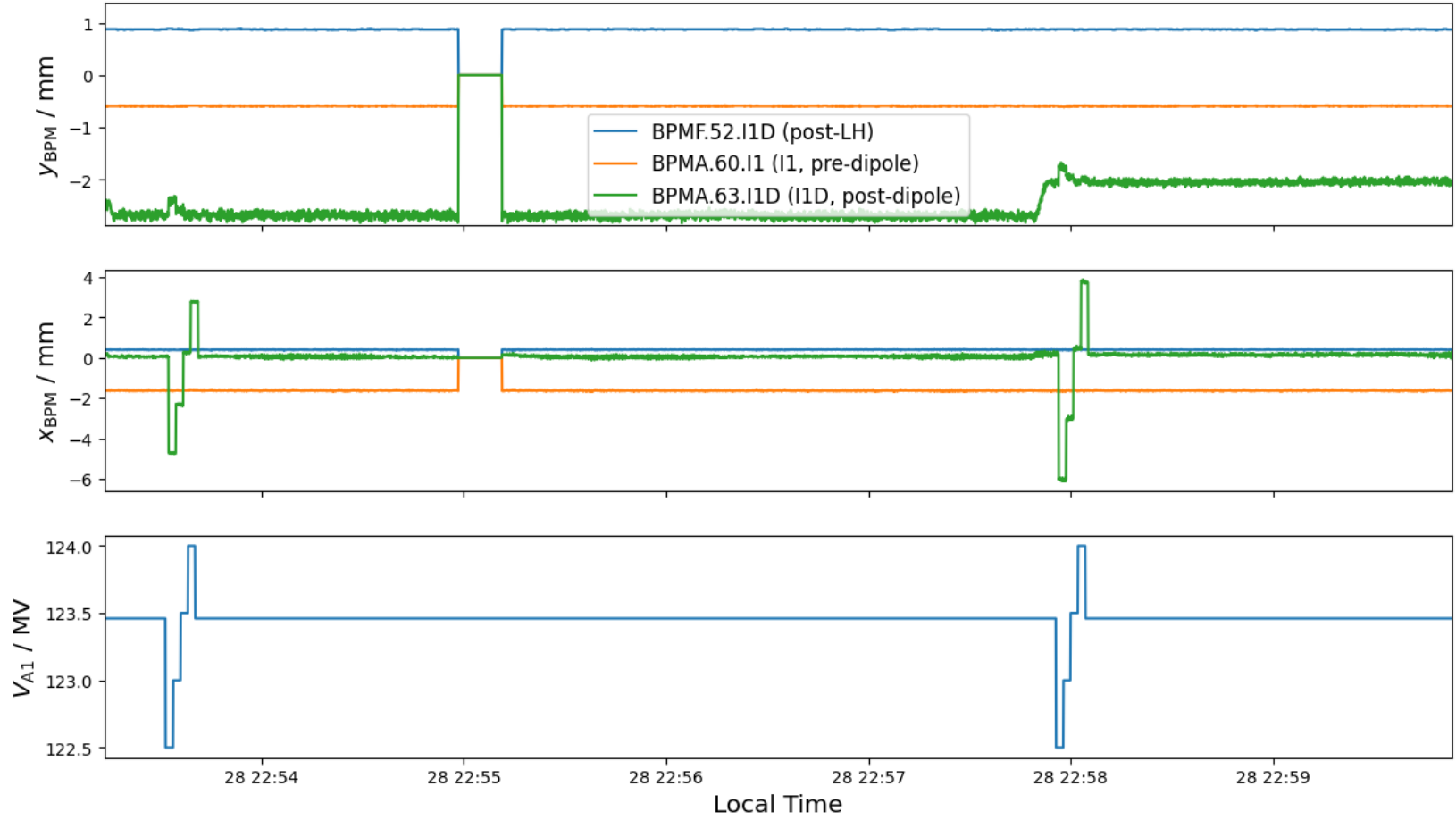
MXL

[stuart.walker@desy.de](mailto:stuart.walker@desy.de)

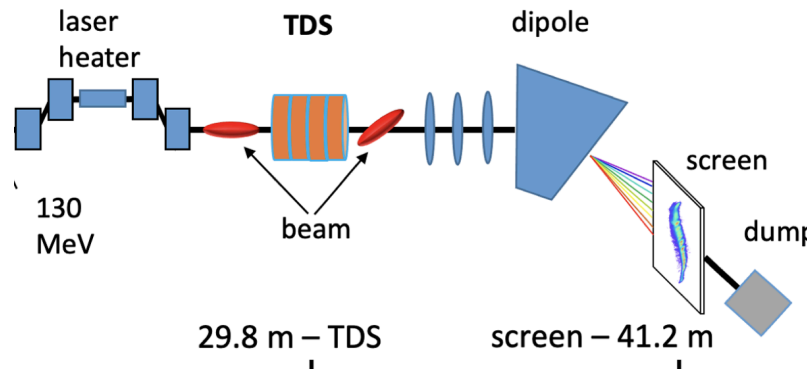
4756

# Backup

# Do we get a kick from A1 when measuring the dispersion?



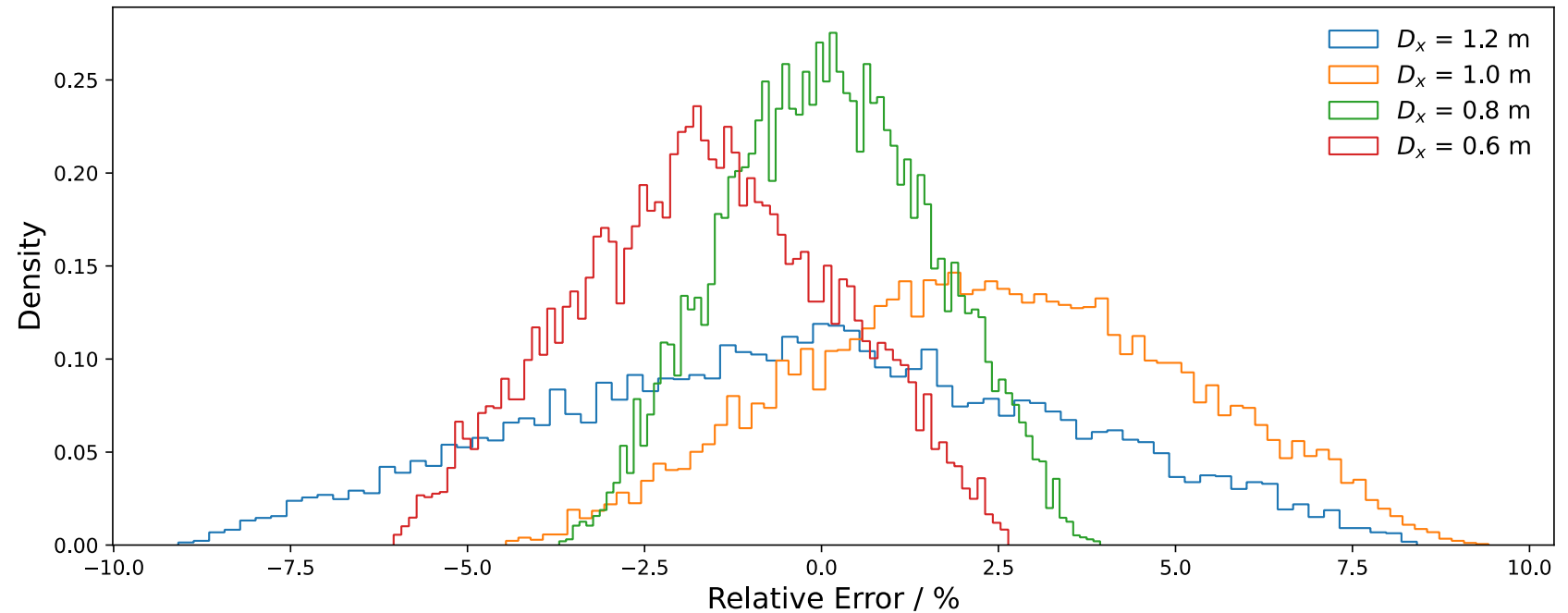
Measurement at lowest dispersion (~0.6m) followed by at highest dispersion (1.1m).



# Possible Dispersion Leakage from the LH Chicane

- We did not remeasure the dispersion at each setpoint of the laser heater  $R_{56}$  scan.
  - Possible: we have some leaked dispersion from the chicane at the screen.
- First two dipoles have same power supplies, second two are each independent.
- Consider: Identical 15% error in first two dipoles and independent 15% errors in second pair, perhaps some sort of "worst case".

- Simulated in OCELOT.
- Dispersion at the screen is tolerant to even large changes in the LH dipoles.
- Nevertheless should measure the dispersion in future experiments.
- Translates max  $\pm 10\%$  error in energy spread.



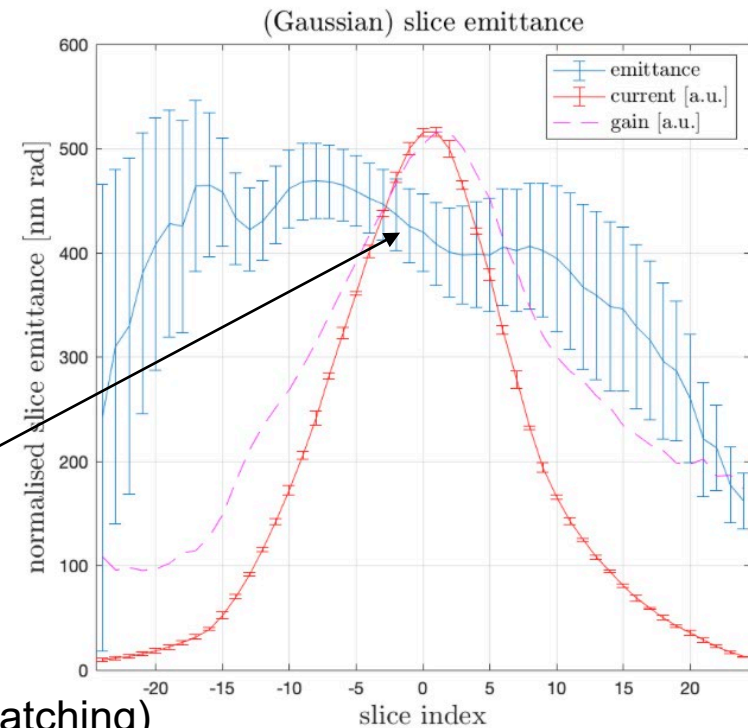
# Emittance Measurements

- Matthias matching/measurement tool is vital to our measurement.
- We match projected and then measure slice—and normally this gives reasonable matching, but we extract much smaller emittances from our scans.
- We see measured central slice emittances that are larger than the corresponding projected emittances even with apparently "good" matching.
- Is there any way we can reduce the error bars here?

```
01.09.2023 00:19 Other
quad scan, reference position QI.52.I1.
Data saved in /home/xfeloper/data/quad_scan
Beam Energy: 130MeV
measurement device: OTRC.59.I1
used quadrupole magnets: QI.52.I1, QI.53.I1
Twiss parameters and emittance results:
emittance_X emittance_Y BMAG_X BMAG_Y alpha
[mm mrad] [mm mrad]
0.37232 0.45426 1.0451 1.0256 -0.925
```

Projected

Central slice (0.45 mm•mrad)  
measured a minute later (no matching)

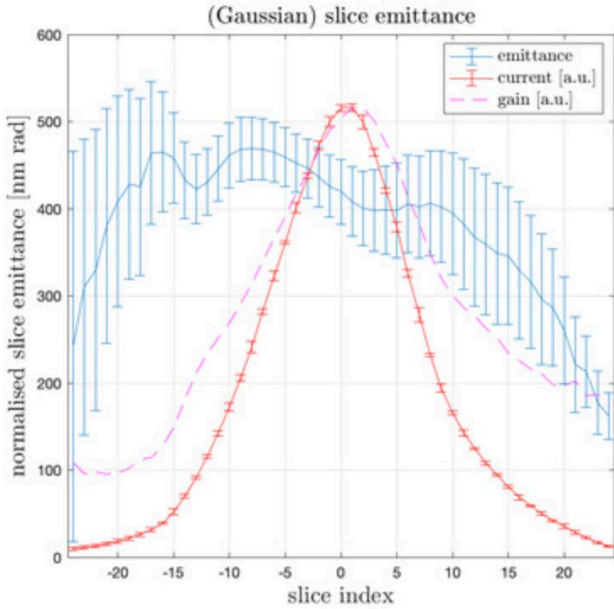
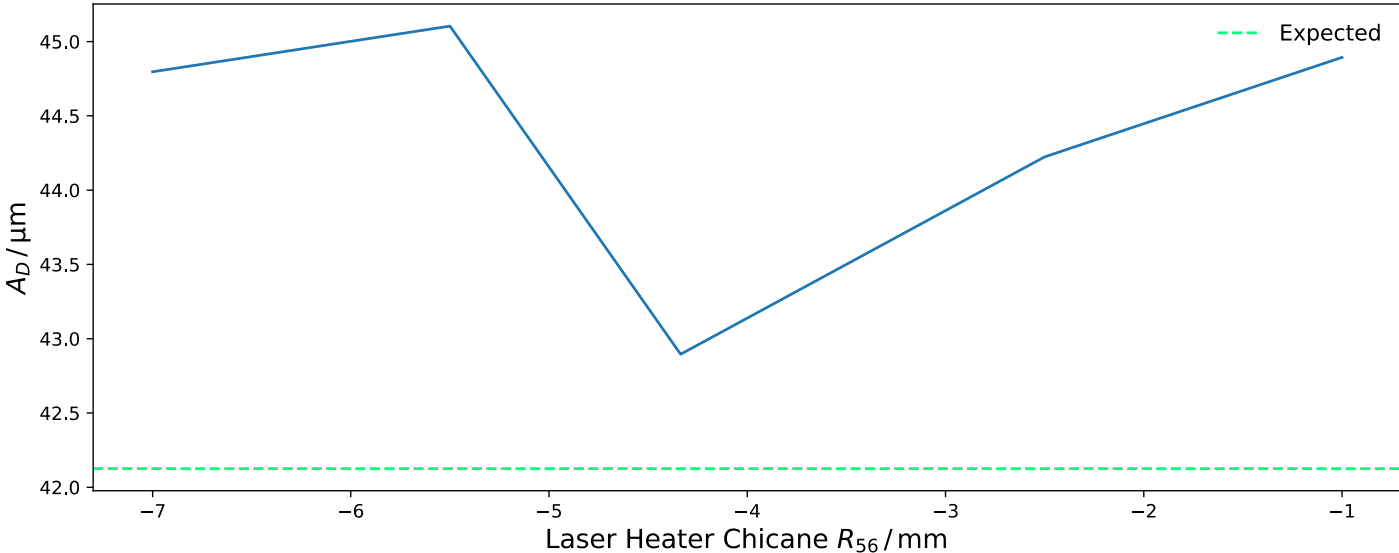


ref

# Constant Contributions to Measured Widths

- In principle both are known:
  - Screen resolution ( $\sigma_R$ ) from dedicated optics simulations (Artem)
  - Betatron contribution ( $\sigma_B$ ) from slice emittance measurement and linear optics (Matthias)
- $A_D = 28\mu\text{m}$  (from Artem and previous scans);  $\beta_x = 0.6\text{ m}$  and  $\epsilon_x = 0.43\text{ mm}\cdot\text{mrad}$  (Matthias matching tool)
- Constant terms remain constant across the scan, but  $\sim 7\%$  larger than expected (from Matthias+Artem).

$$A_D = \sigma_R^2 + \sigma_B^2$$

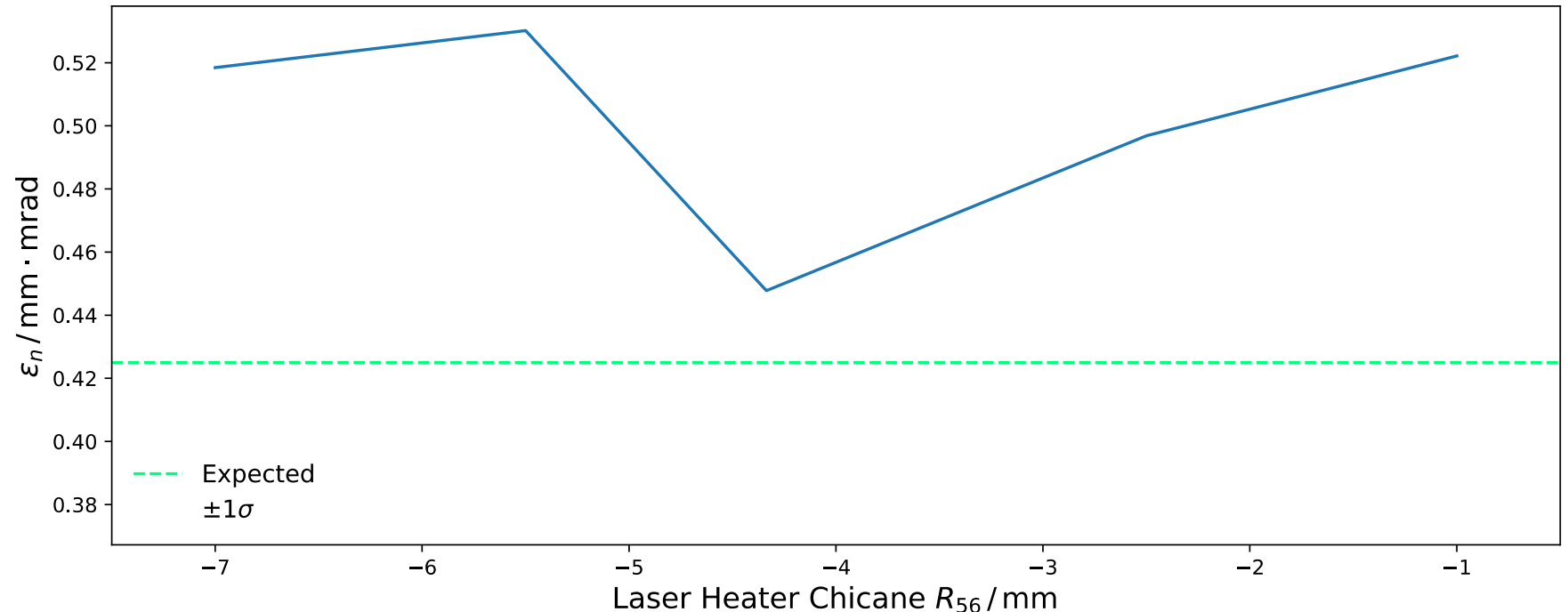




# Deriving the emittance assuming the screen resolution

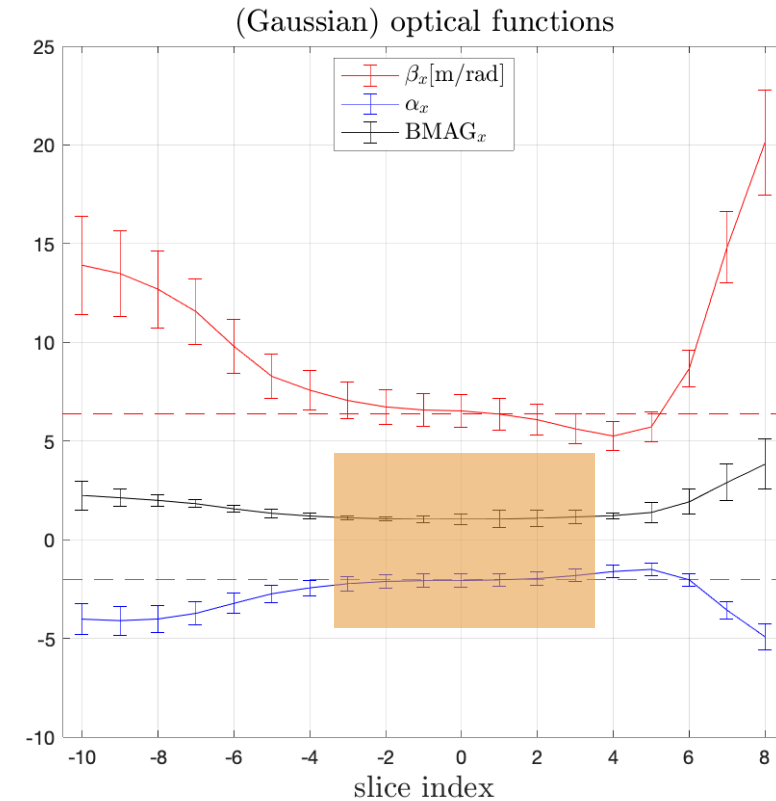
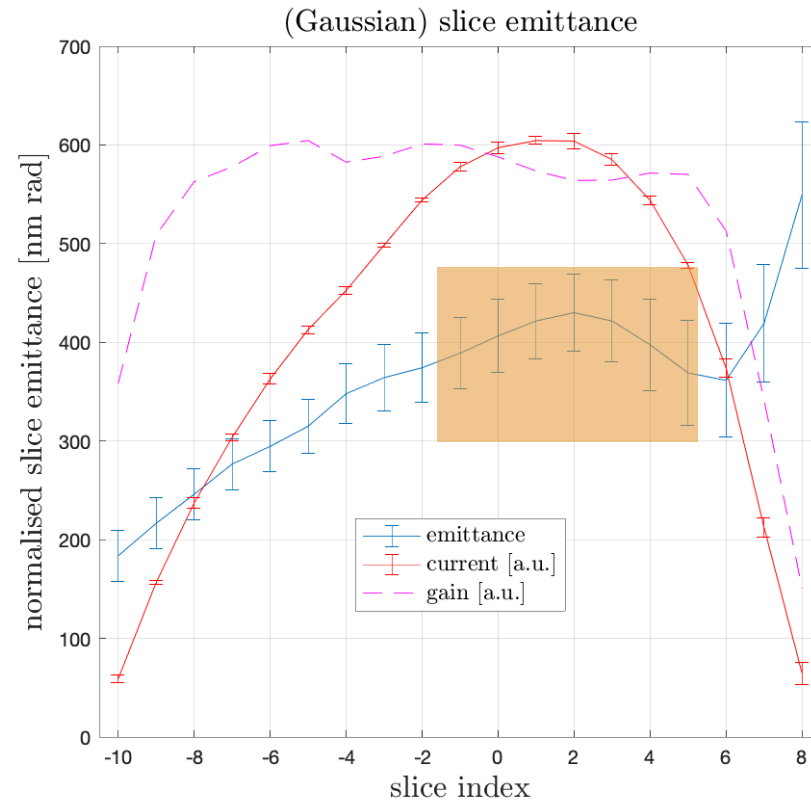
- Fitting  $A_D$  and using a pre-calculated value of  $\sigma_R$ , we can extract a value for the slice emittance.
- We independently measure, using Matthias' matching tool,  $(0.43 \pm 0.05)$  mm•mrad for the slice emittance.
- Not entirely in disagreement with our derived value.
- Can we reduce the uncertainties in the emittance measurement?

$$\epsilon_n = \frac{\gamma_0}{\beta_x} (A_D - \sigma_R^2)$$



# Beam Matching

- Central slice always matched for measurements
- Very important, otherwise can get an artificially deflated energy spread result!
- Use screen just in front of TDS
- We remeasured the dispersion at every single point in the dispersion scan and once per TDS scan.



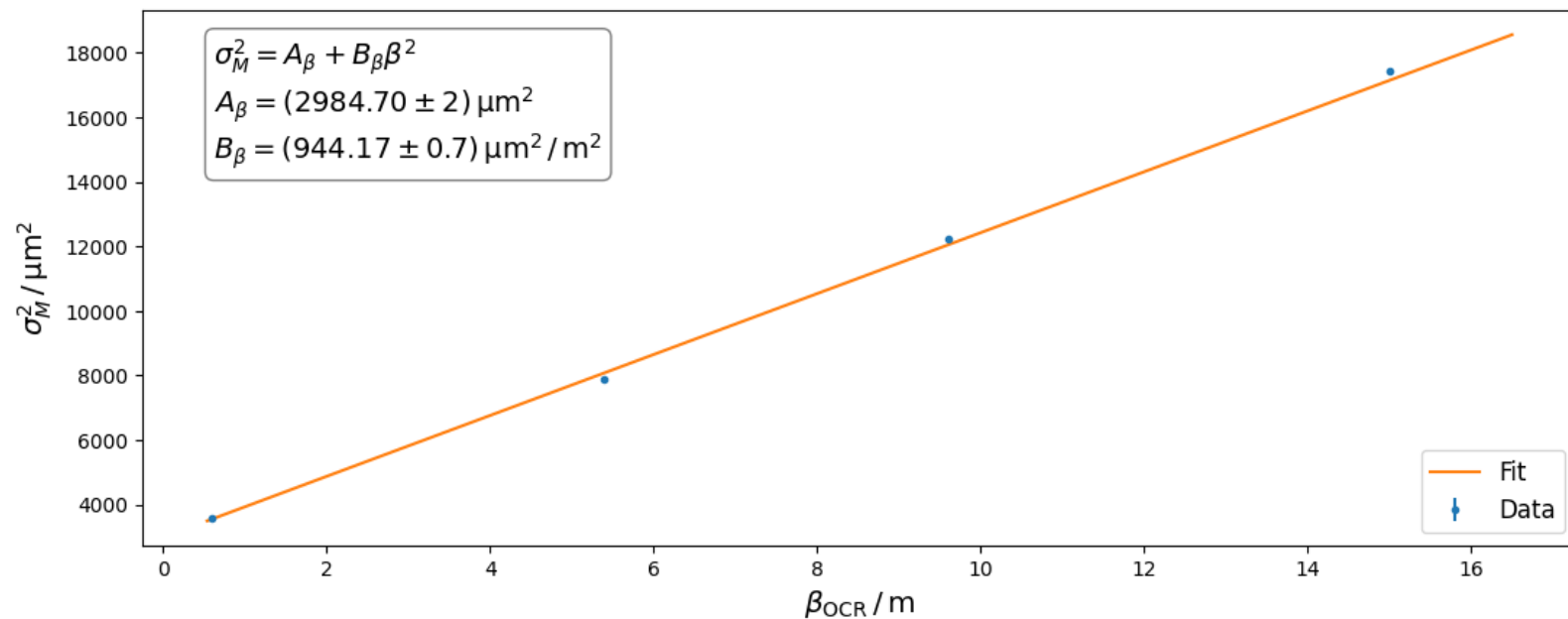
# Beta Scan

$$\sigma_M^2 = \sigma_R^2 + \sigma_B^2 + \frac{D^2}{E^2} \sigma_E^2 + \frac{D^2}{E^2} (ekV)^2 \sigma_x^2$$

Vary

Fix

$$\sigma_M^2 = A_\beta + \sigma_B^2 = A_\beta + B_\beta \beta_x^2$$



$$\sigma_M^2 = \sigma_R^2 + \sigma_B^2 + \frac{D^2}{E^2} \sigma_E^2 + \frac{D^2}{E^2} (ekV)^2 \sigma_I^2$$

$$\sigma_M^2 = A_V + B_V V^2$$

$$\sigma_M^2 = A_D + B_D D^2$$

$$\sigma_M^2 = A_\beta + B_\beta \beta_x$$

$$\sigma_E = \frac{E_0}{D_0} \sqrt{A_V - A_D}$$

$$\sigma_I = \frac{E_0}{D_0 ek} \sqrt{B_V}$$

$$\sigma_B = \sqrt{B_\beta \beta_x^0}$$

$$B_\beta = \sigma_I^2 (\beta_y^0 + 0.25 L^2 \gamma_y^0 - L \alpha_y^0)^{-1}$$

$$\sigma_R = \sqrt{A_D - \sigma_B^2}$$

$$\sigma_E = \frac{E_0}{D_0} \sqrt{D_0^2 B_D - V_0^2 B_V}$$

$$\sigma_I = \frac{E_0}{D_0 ek V_0} \sqrt{D_0^2 B_D + A_D - A_V}$$

$$\sigma_B = \sqrt{A_D + B_D D_0^2 - A_\beta}$$

$$B_\beta = \varepsilon_n \gamma_0^{-1}$$

$$\sigma_R = \sqrt{A_\beta - B_D D_0^2}$$

$$\sigma_M^2 = \sigma_R^2 + \sigma_B^2 + \frac{D^2}{E^2} \sigma_E^2 + \frac{D^2}{E^2} (ekV)^2 \sigma_I^2$$

Variable	Value	Alt. Value	Units
$V_0$	-0.63	-	MV
$D_0$	1.169	-	m
$A_V$	$(3.043 \pm 0.002) \times 10^{-9}$	-	$m^2$
$B_V$	$(1.326 \pm 0.003) \times 10^{-21}$	-	$m^2 V^{-2}$
$A_D$	$(1.540 \pm 0.002) \times 10^{-9}$	-	$m^2$
$B_D$	$(1.416 \pm 0.002) \times 10^{-9}$	-	
$A_\beta$	-	$(2.985 \pm 0.002) \times 10^{-9}$	$m^2$
$B_\beta$	-	$(9.442 \pm 0.007) \times 10^{-10}$	m
$\sigma_E$	$4.313 \pm 0.004$	$4.187 \pm 0.005$	keV
$\sigma_I$	$(6.444 \pm 0.008) \times 10^{-5}$	$(5.88 \pm 0.03) \times 10^{-5}$	m
$\sigma_B$	$(2.834 \pm 0.004) \times 10^{-5}$	$(2.215 \pm 0.009) \times 10^{-5}$	m
$\sigma_R$	$(2.714 \pm 0.005) \times 10^{-5}$	$(3.239 \pm 0.005) \times 10^{-5}$	m
$\varepsilon_x$	$0.3408 \pm 0.0009$	$0.2403 \pm 0.0002$	mm · mrad

2022 data with beta scan  
(Attempting to recreate 2021)

# New Tooling

## TDS Calibration

- Additionally the TDS needs to be calibrated for each measurement campaign to get the voltage → also takes up to an hour.
- Also developed an injector TDS calibration GUI
  - Automatically scan TDS “amplitudes” and phases to build mapping of amplitudes [%] to voltages [MV] → Gives all downstream longitudinal calibrations!
  - Also useful for non-invasive emittance and bunch length measurements.
- **Bottom line:** calibrate the injector TDS in a few minutes with no human intervention.

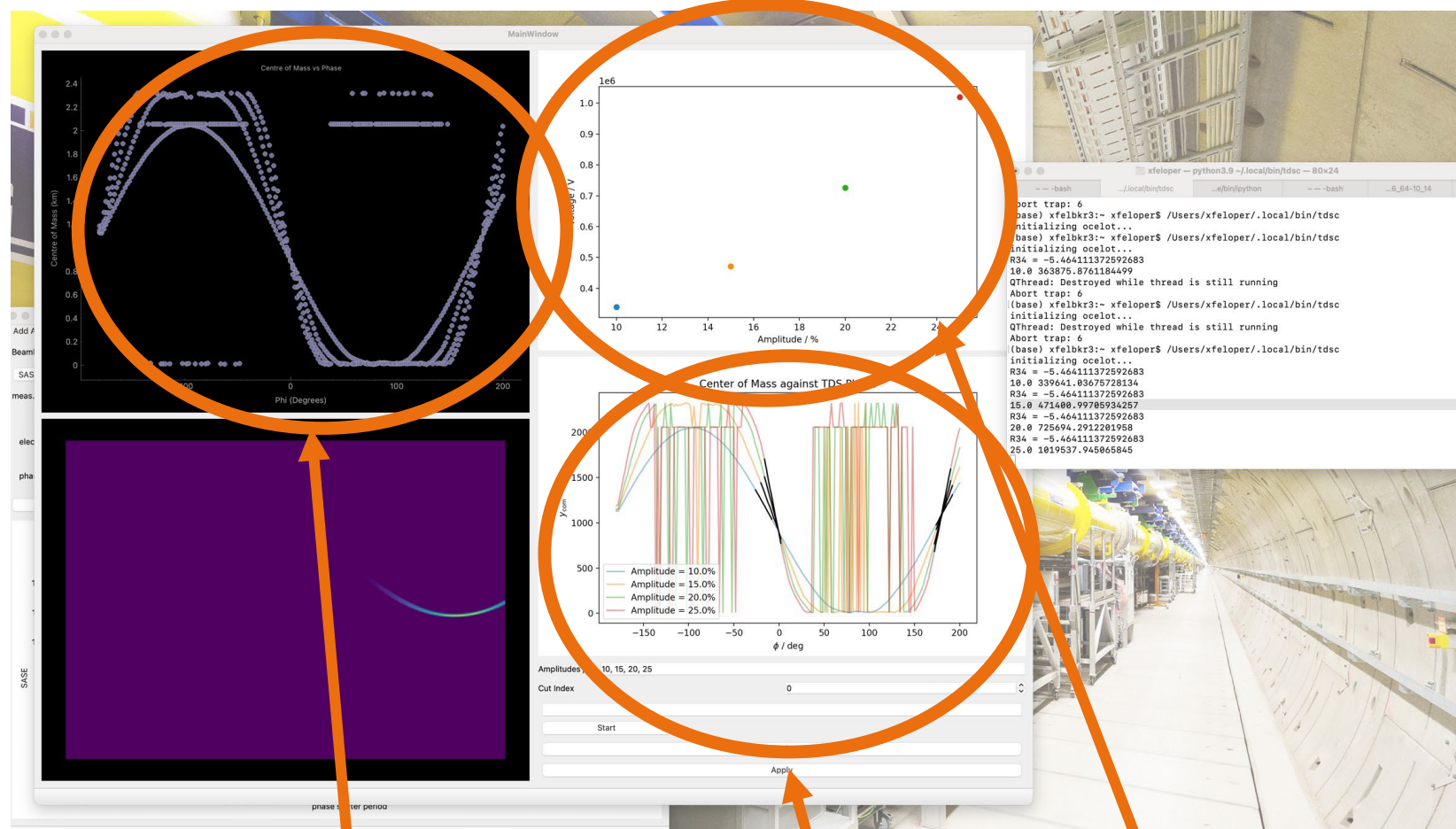


Image centre of mass vs phase

B2 calibrator development ongoing.

Non-invasive development ongoing.

Amplitude to voltage mapping (the calibration final result).

Extracted zero-crossing gradients at each amplitude